




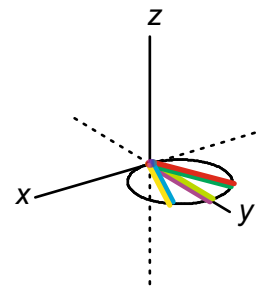
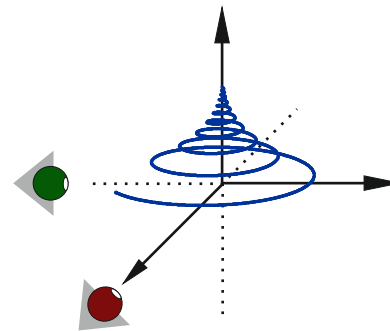


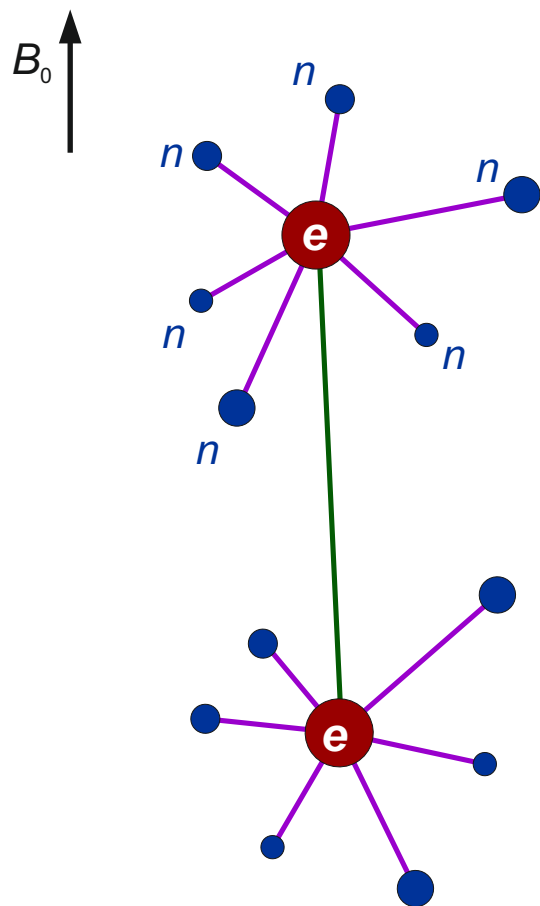
# Pulsed EPR spectroscopy with a focus on DEER (PELDOR)

Gunnar Jeschke  
ETH Zürich, Lab. Phys. Chem.

-  Why pulsed EPR?
-  How does it work? FID & echoes
-  A few concepts and pulse sequences
-  DEER: Basics and implementation
-  DEER: Data analysis

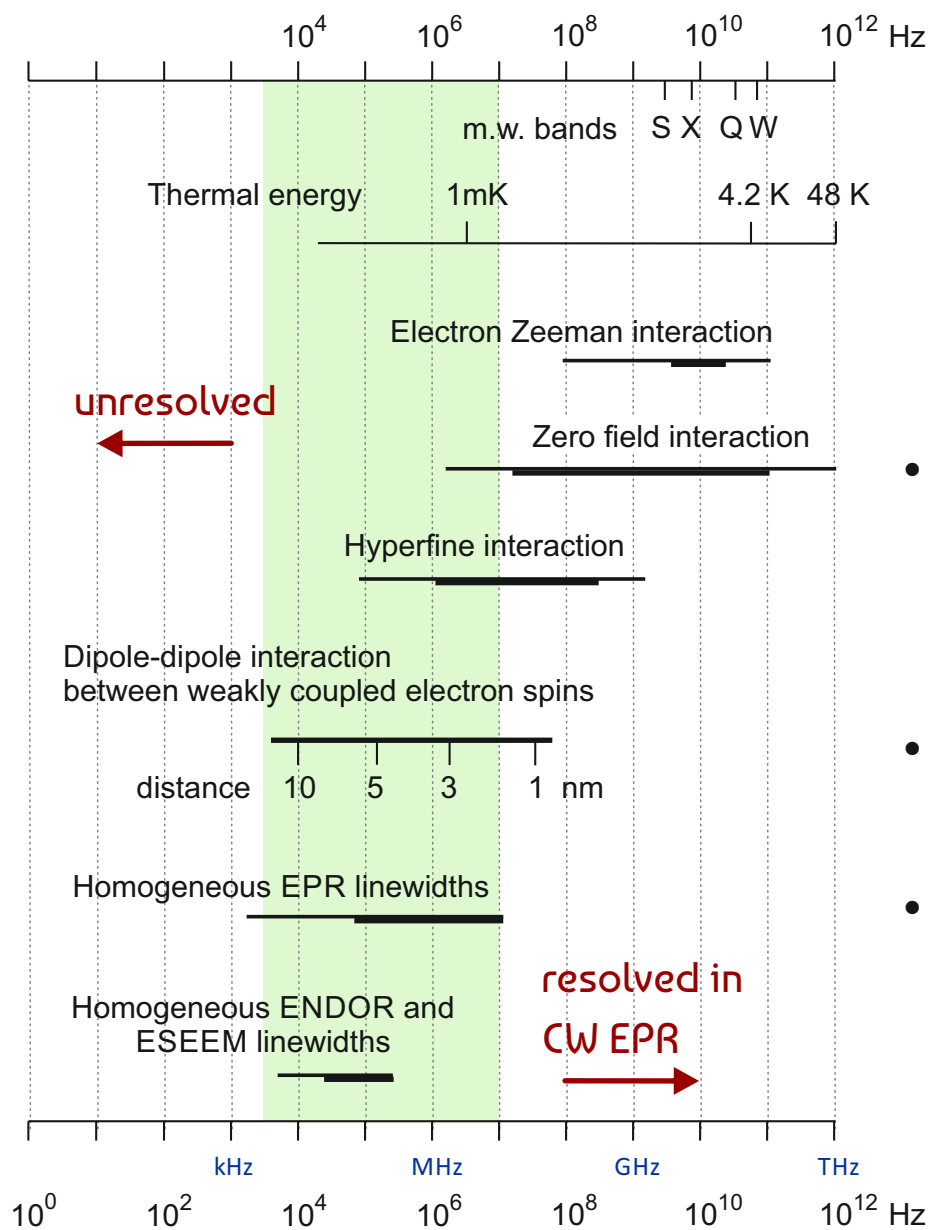


## Interactions and the information that they provide



Name	Information
electron Zeeman	fingerprinting of radical type or metal coordination
hyperfine	distribution of the SOMO (reactivity) distance of protons from the center of spin density
nuclear Zeeman	identification of nuclei that give rise to hfi
nuclear quadrupole	binding situation of the nucleus for $I > 1/2$ (chemical shift is not available)
zero-field	fingerprinting of triplet type or metal coordination spin state for metal ions (low or high spin)
exchange	orbital overlap (important for electron transfer)
dipole-dipole	distances in the nanometer range (15 - 100 Å)

## An overview of microwave bands and interactions



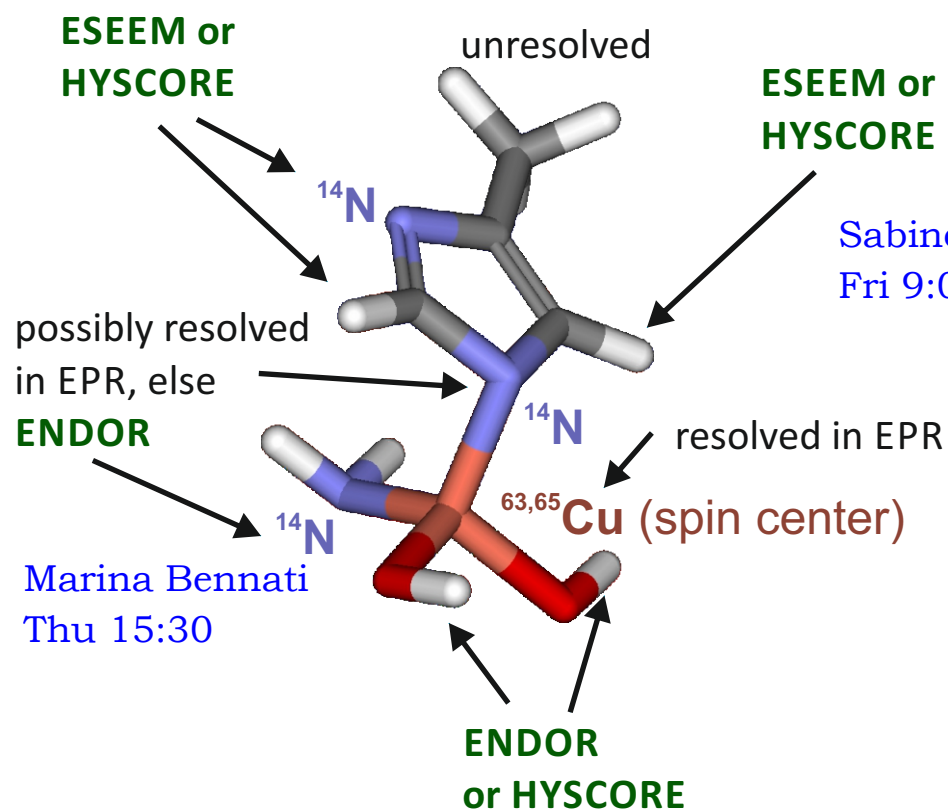
By separating interactions, pulsed EPR provides information lost by line broadening in CW EPR

- for electron group spin  $> 1/2$  (more than one unpaired electron)
- most valuable source of EPR restraints on structure
- resolution limit depends on sample preparation

1 mT corresponds to 28 MHz

## Example: Measuring hyperfine couplings

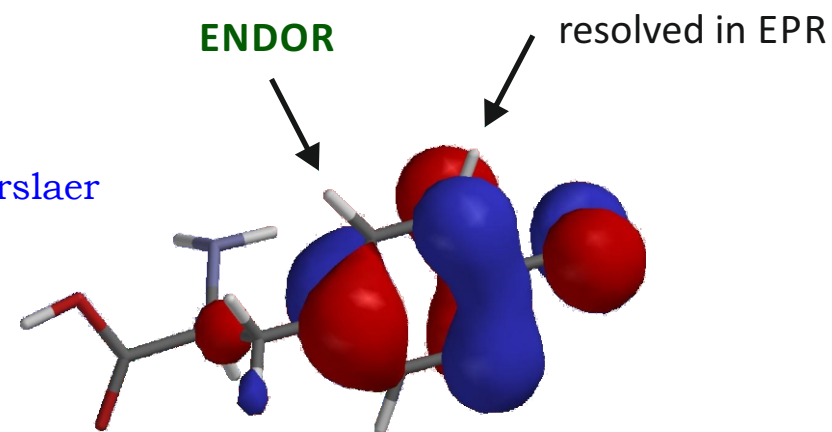
Histidine-coordinated Cu(II)



Sabine Van Doorslaer  
Fri 9:00

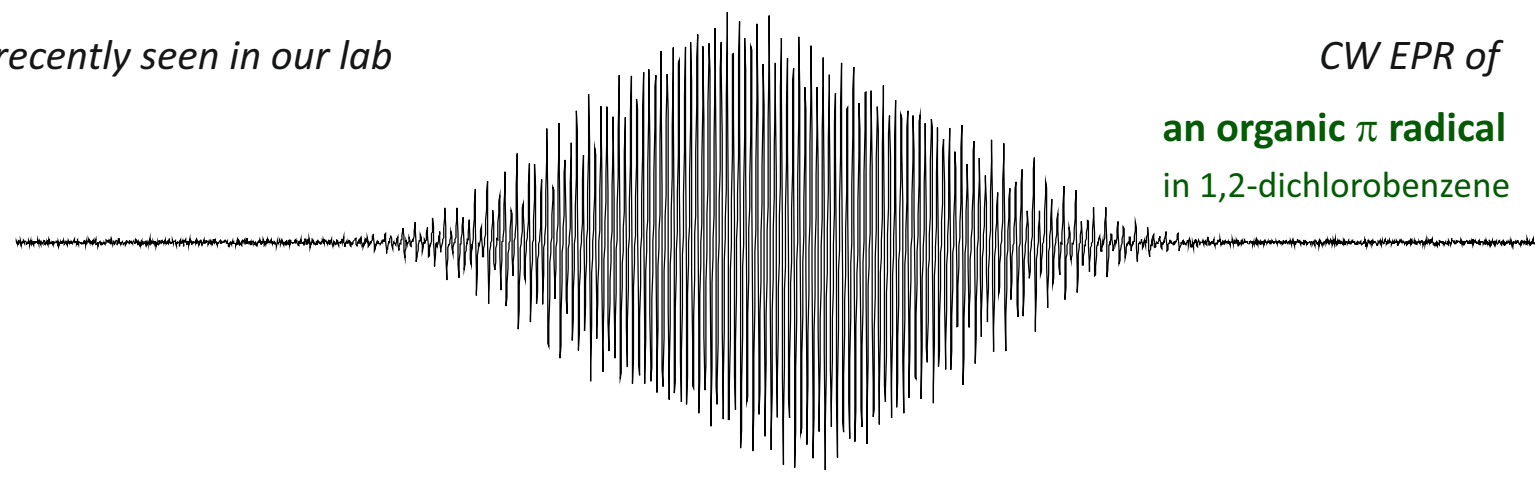
Marina Bennati  
Thu 15:30

Tyrosyl radical

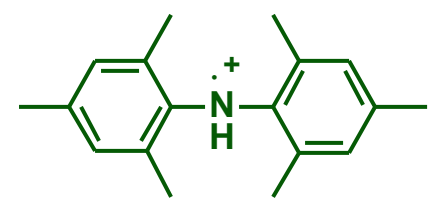


## Even if it's resolved, ...

recently seen in our lab



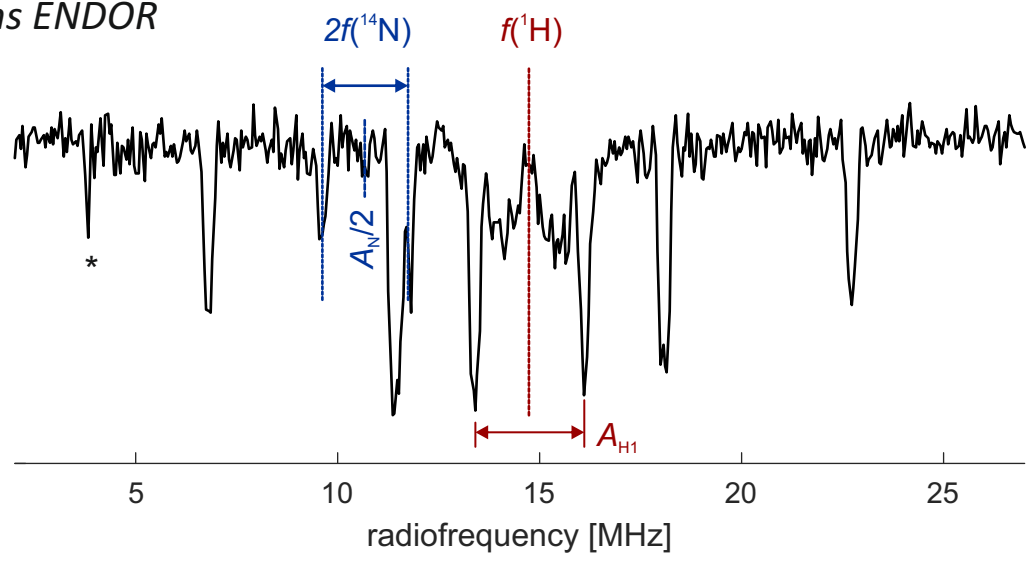
CW EPR of  
an organic  $\pi$  radical  
in 1,2-dichlorobenzene



sample courtesy  
Agnes Kütt  
University Tallin

## a more complicated experiment can make it simpler

Mims ENDOR



- ⊕ identity of the nuclei ( $^1\text{H}$ ,  $^{14}\text{N}$ )
- ⊕ direct read-off of hyperfine couplings

# Pulsed EPR is a zoo of techniques

## EPR spectra

- Fourier-transform EPR
- field-swept echo-detected EPR

## Electron spin relaxation

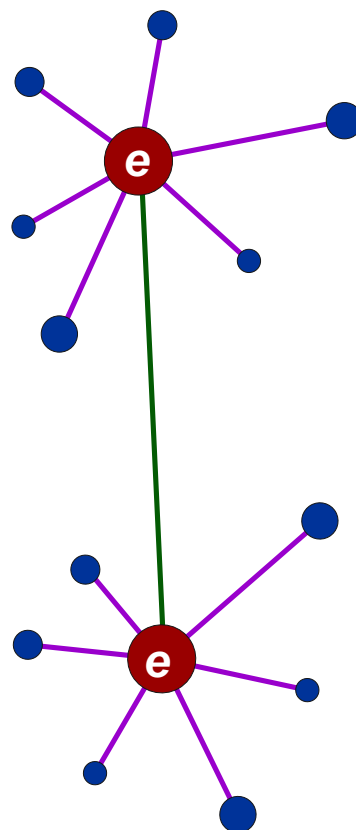
- Hahn echo decay ( $T_2/T_m$ )
- Inversion recovery ( $T_1$ )

## Hyperfine couplings nuclear Zeeman & nuclear quadrupole interactions

- ESEEM Sabine, Fri 9:00
- HYSORE
- Davies ENDOR
- Mims ENDOR Marina, Thu 15:30
- ELDOR-detected NMR Daniella, Fri 11:00

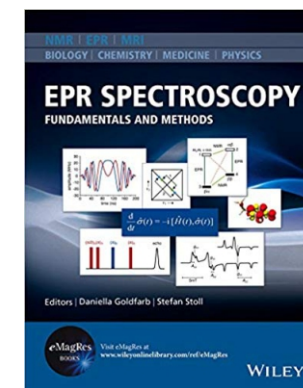
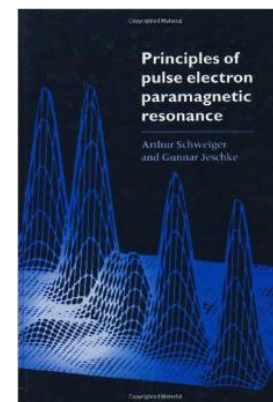
## Electron-electron couplings

- DEER/PELDOR
- RIDME
- SIFTER
- DQC-EPR



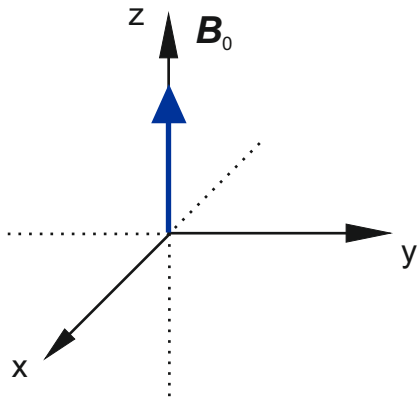
### Fine print

- more techniques exist
- new ones are still coming up
- almost all of them fit into one of the four boxes

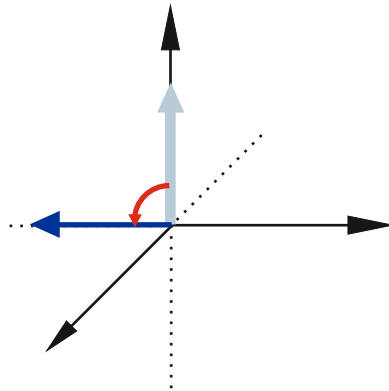


# Free induction decay (FID) and Fourier transform EPR

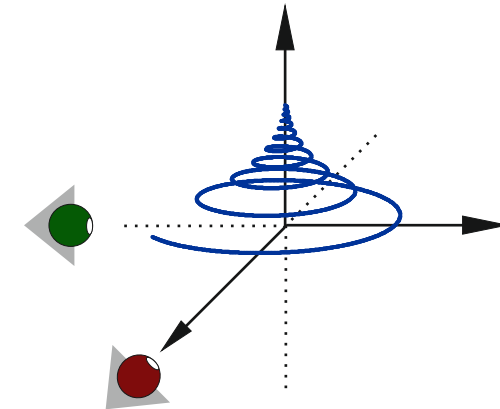
Thermal equilibrium



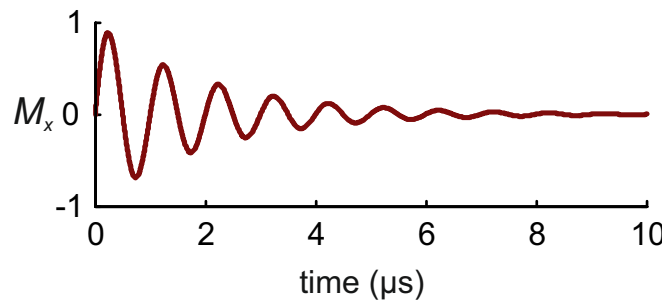
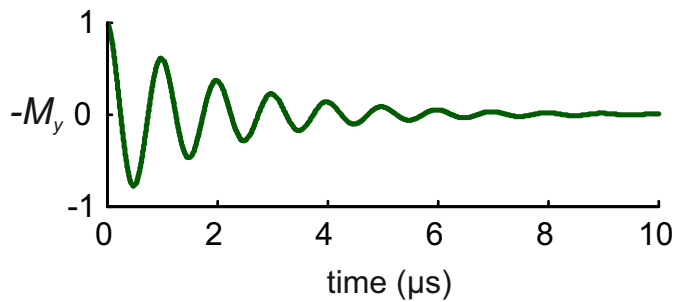
$\pi/2$  ( $90^\circ$ ) pulse



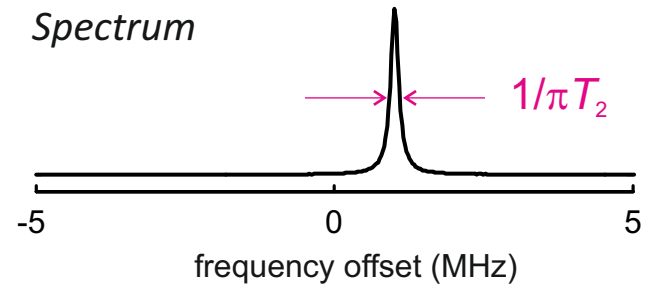
Free evolution (FID)



Complex (quadrature) signal:  $-M_y + i M_x$



Fourier transform



# Motion of the magnetization vector without relaxation

## Classical equation of motion

spin has, both, angular momentum  $J$  and a magnetic moment  $\mu$

$$\frac{d\vec{J}}{dt} = \underbrace{\vec{\mu} \times \vec{B}(t)}_{\text{torque}} \quad \vec{J} \text{ total angular momentum vector, } \vec{B}(t) \text{ includes static and oscillatory magnetic fields}$$

as  $\gamma \vec{J} = \vec{\mu}$  and  $\vec{M} = \vec{\mu}/V$ , we have

$$\frac{d\vec{M}}{dt} = \vec{M} \times \gamma \vec{B}(t)$$

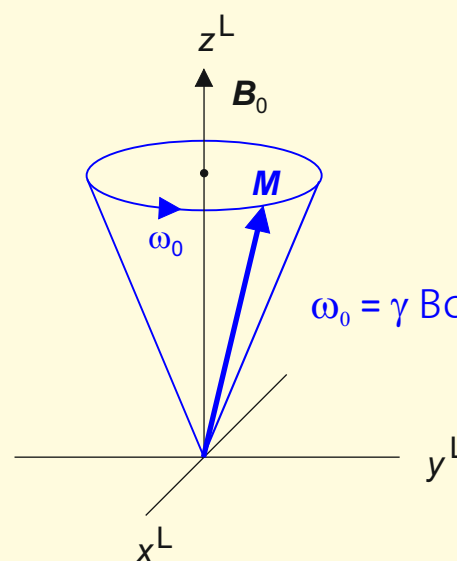
and by expanding the cross product

$$\dot{M}_x = -\gamma B_z(t) M_y + \gamma B_y(t) M_z$$

$$\dot{M}_y = \gamma B_z(t) M_x - \gamma B_x(t) M_z$$

$$\dot{M}_z = \gamma B_x(t) M_y - \gamma B_y(t) M_x$$

*Precession in a static field  $B_0$*



$$\dot{M}_x = -\omega_0 M_y$$

$$\dot{M}_y = \omega_0 M_x$$

$$\dot{M}_z = 0$$



## Relaxation in the magnetization vector model

### Relaxation is due to flips of individual spins

⇒ only phenomenological description possible in a classical picture

⇒ length of  $\vec{M}$  is not a constant of motion

### Two types of relaxation just a reminder, Thomas Prisner talked about this

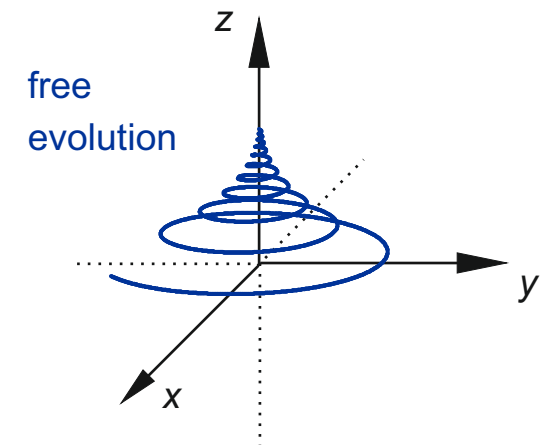
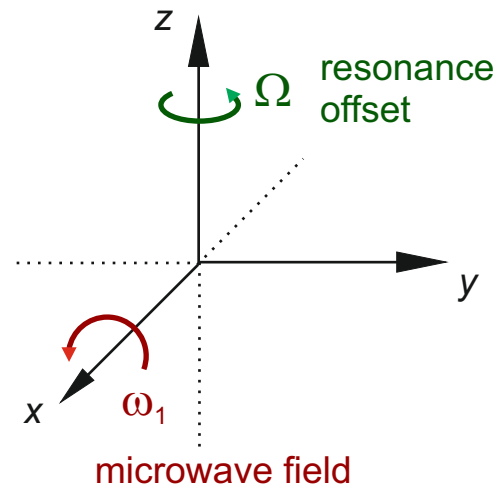
- longitudinal relaxation requires energy exchange with the environment, enthalpic, time constant  $T_1$
- transverse relaxation reduces phase coherence of the ensemble, entropic, time constant  $T_2$

### Bloch equations with relaxation rotating frame

$$\dot{M}_x = -\Omega M_y - \frac{M_x}{T_2}$$

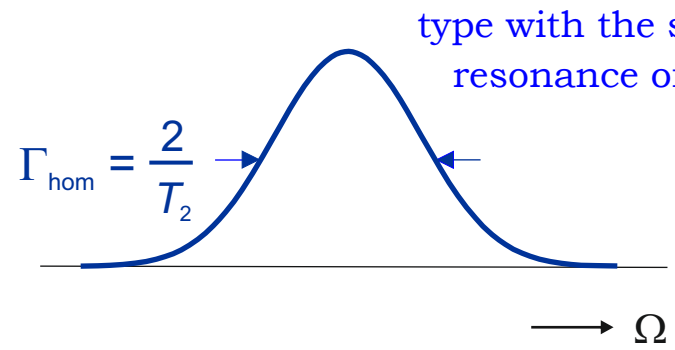
$$\dot{M}_y = \Omega M_x - \omega_1 M_z - \frac{M_y}{T_2}$$

$$\dot{M}_z = \omega_1 M_y - \frac{M_z - M_0}{T_1}$$



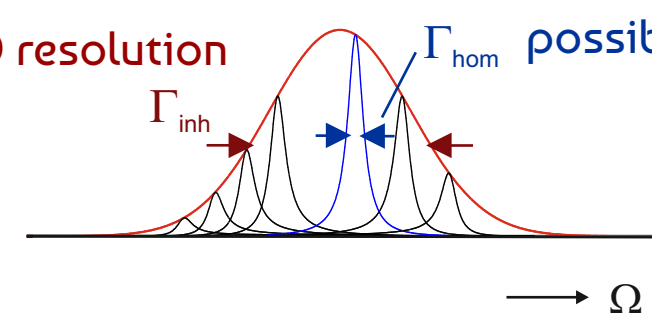
# Homogeneous and inhomogeneous lines

**Homogeneous line** all spins of the same type with the same resonance offset

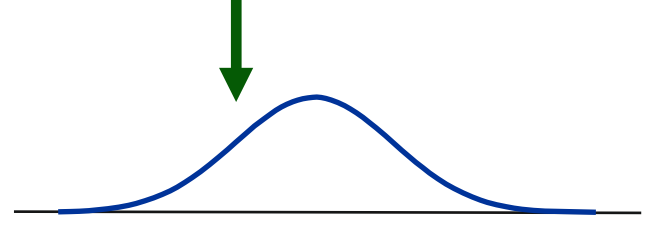


**Inhomogeneous line**

**FID resolution**  $\Gamma_{\text{inh}}$  **possible resolution**  $\Gamma_{\text{hom}}$

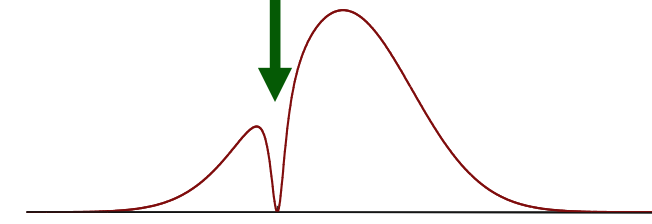


**Irradiation** (long, weak pulse)



- saturation

**Irradiation** (long, weak pulse)

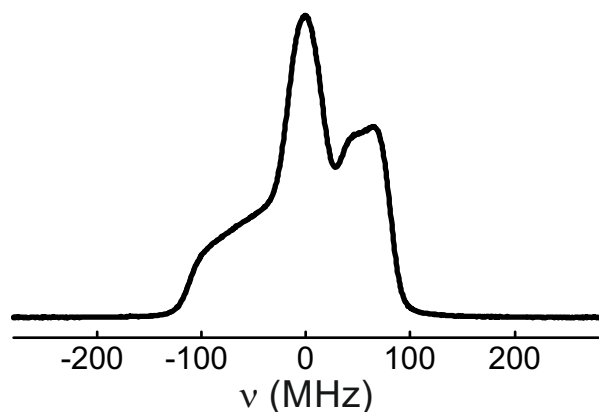


- only a spin packet is saturated
- spectral hole burning  
(used in Davies ENDOR, ELDOR-detected EPR)

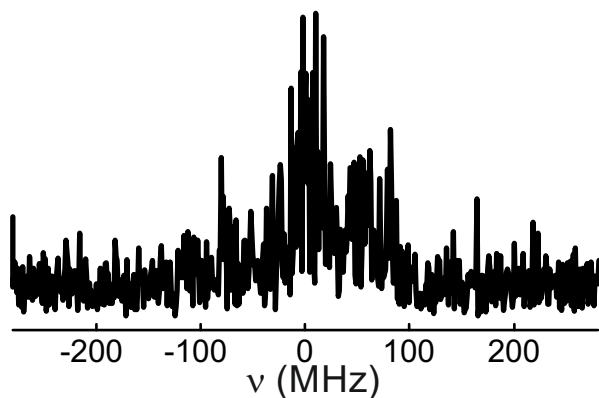
**!** in the solid state, inhomogeneous broadening by anisotropic interactions is often dominant

# Why a single pulse is not sufficient for solids

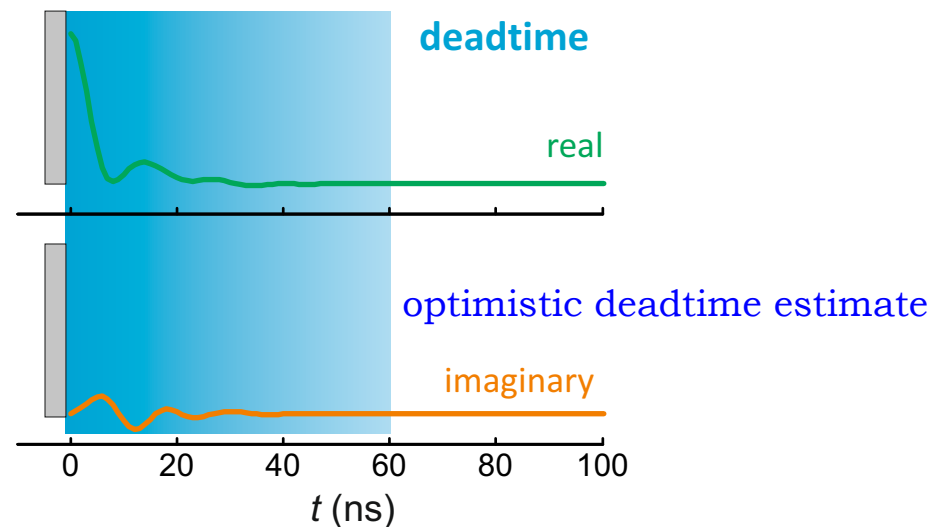
EPR absorption spectrum of a nitroxide (X band)



FT of the accessible part of the FID



... and the corresponding FID



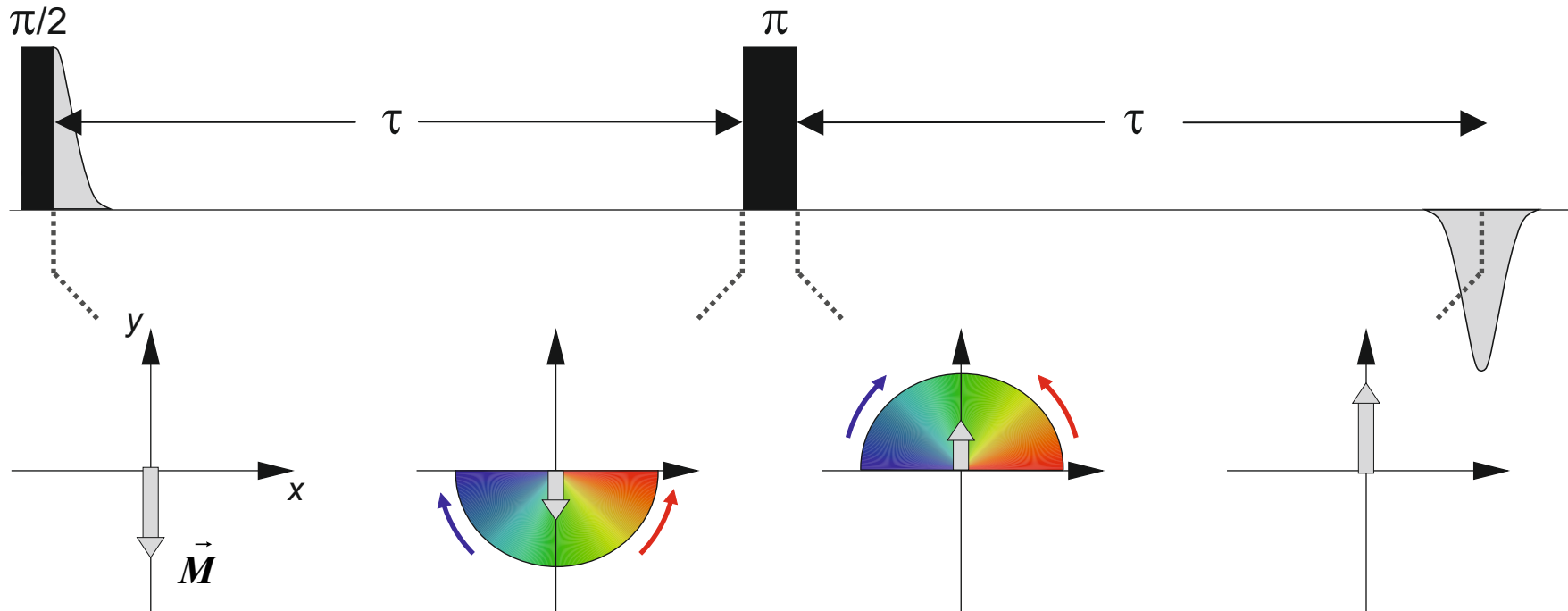
- destructive interference of spin packets with different resonance offset dampens FID within dead time

FT EPR by FID detection possible only for:

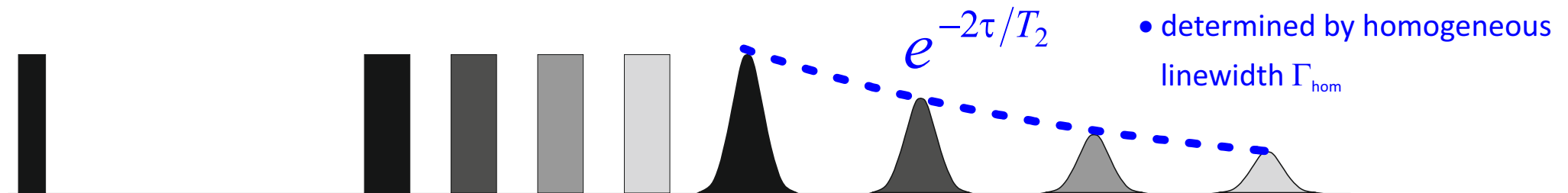
- spectral widths < 100 MHz\*
- line widths  $\approx 3$  MHz

\*with modern arbitrary wavefunction generators: up to 800 MHz

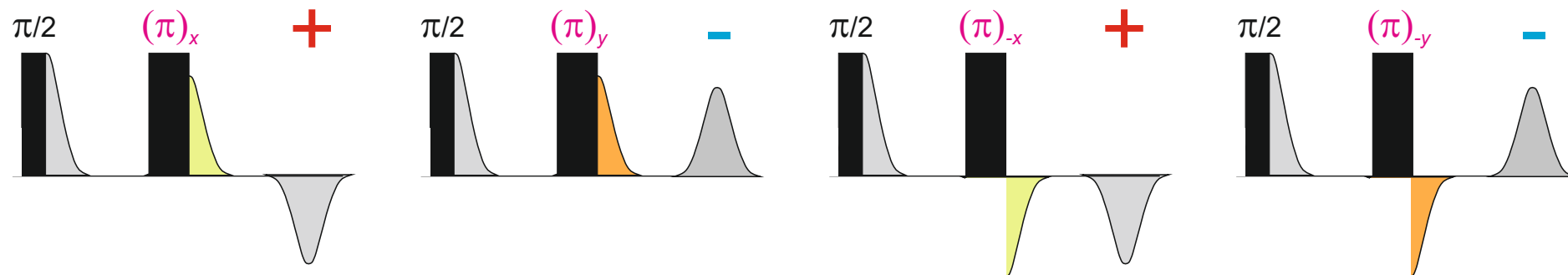
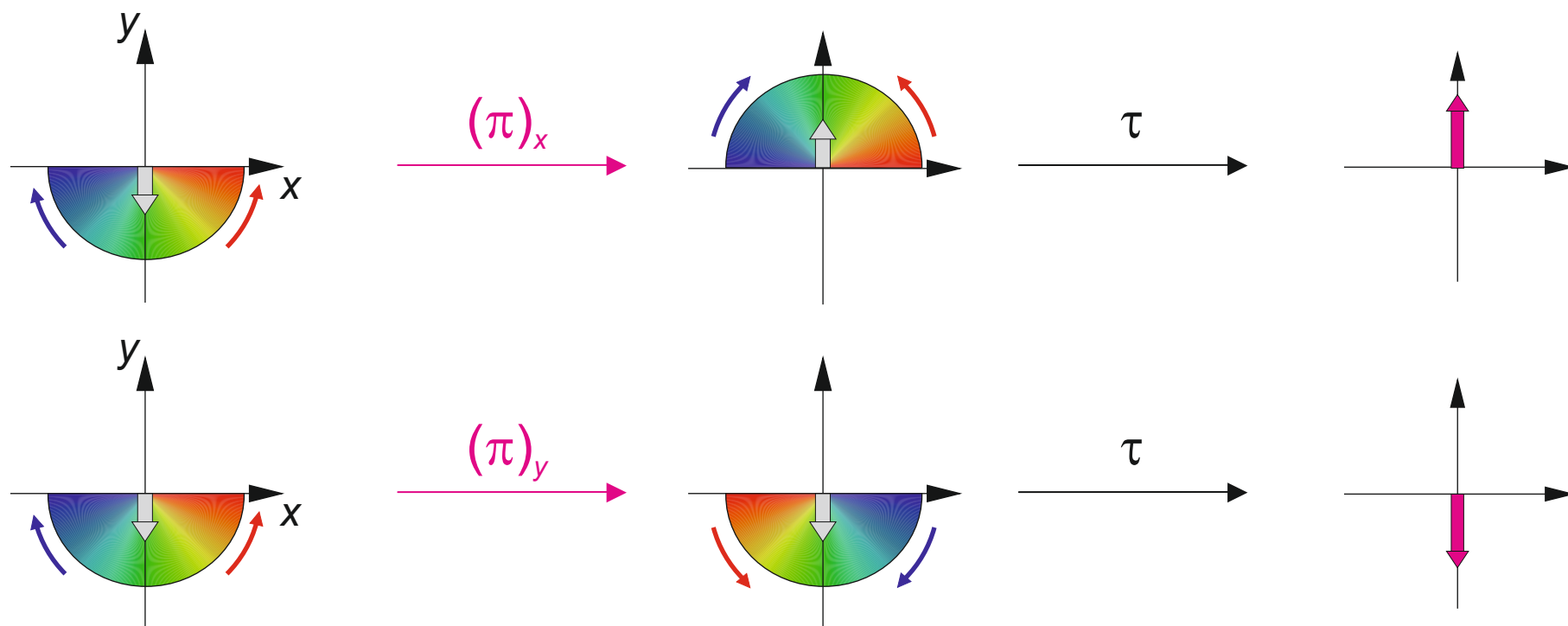
# The primary echo (Hahn echo)



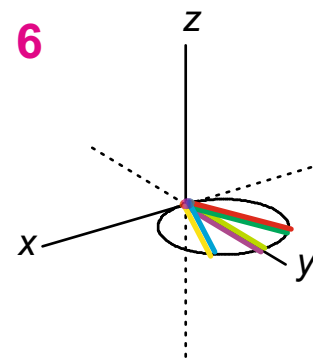
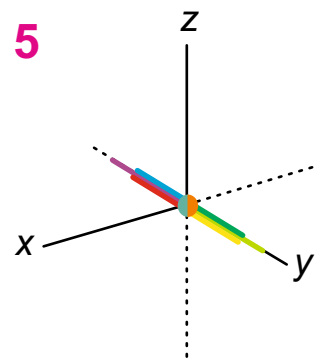
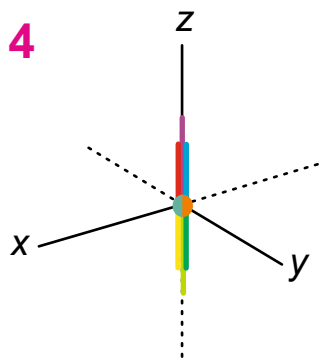
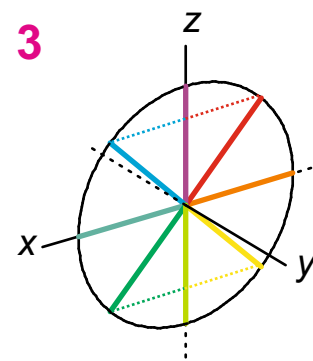
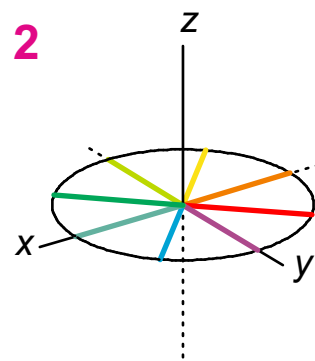
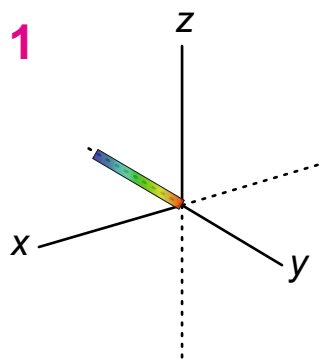
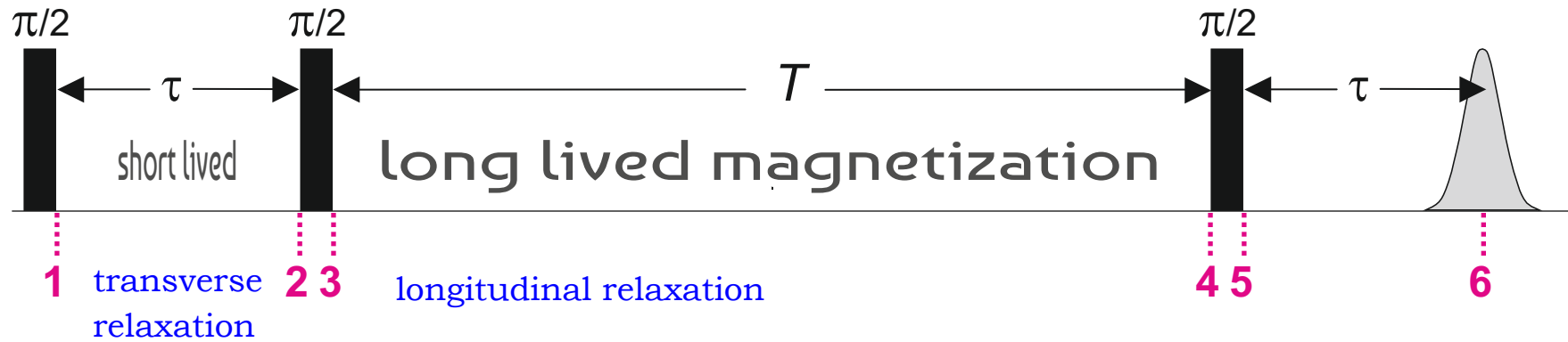
## Echo decay



# Basics of phase cycling

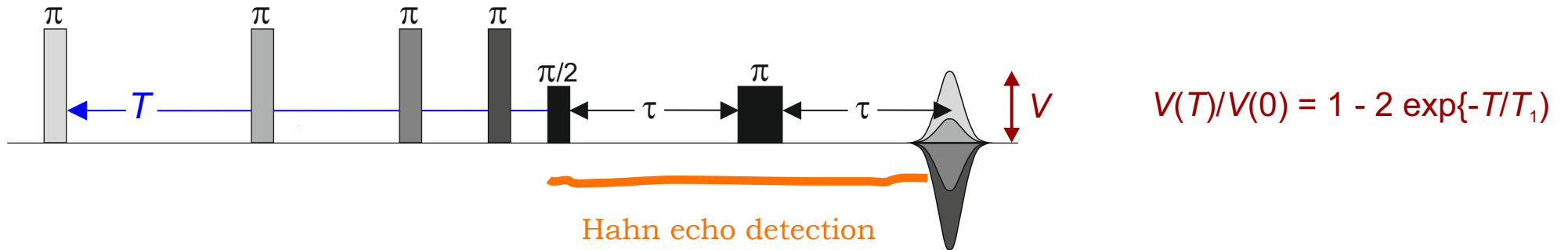


# The stimulated echo experiment



- half of the magnetization is lost by relaxation
- the other half decays *approximately* with  $1/T_1$  during time  $T$
- typically, some kind of magnetization transfer during time  $T$  is observed

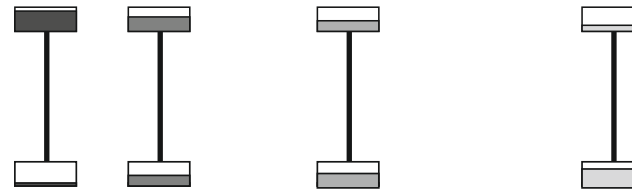
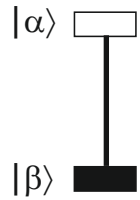
# The inversion recovery pulse sequence



Thermal equilibrium

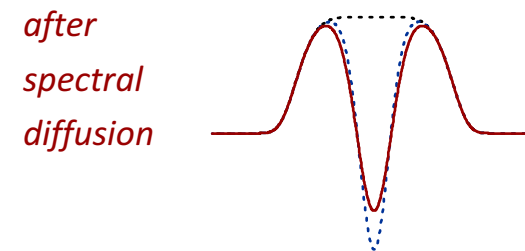
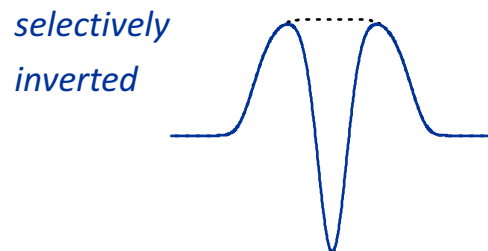
After inversion

Longitudinal relaxation



$T \rightarrow$

**Beware of spectral diffusion** spins change resonance frequency or exchange magnetization



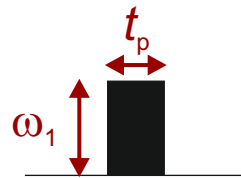
# Nutation and excitation bandwidth

Rotating-frame Hamiltonian during irradiation

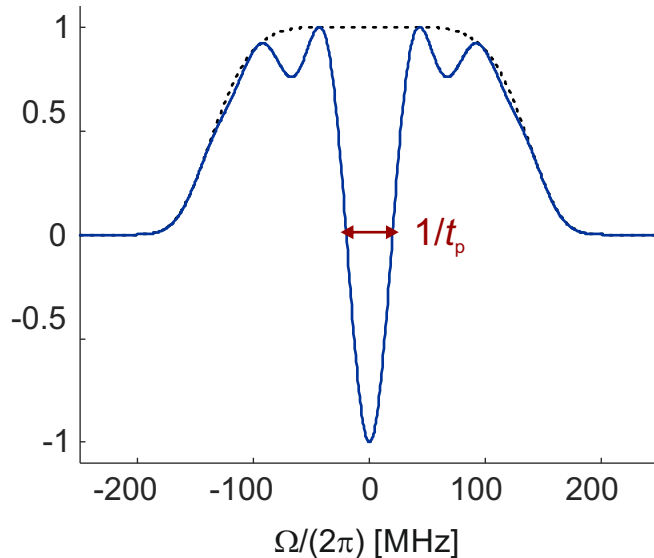
$$\mathcal{H} = \Omega S_z + \omega_1 S_x$$

The flip angle is defined on resonance

$$\beta = \omega_1 t_p$$



Excitation profile



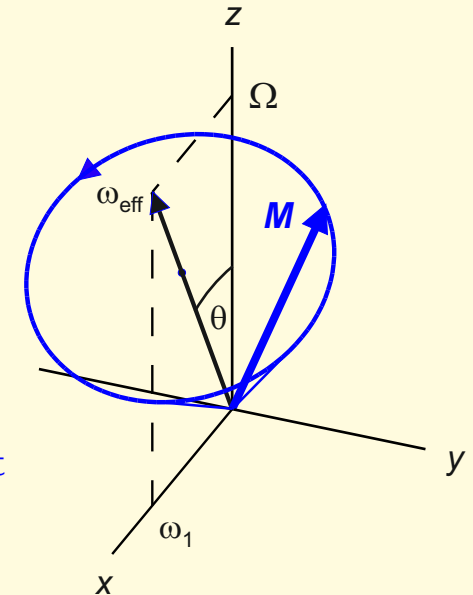
Simulation for  
 $\beta = \pi$   
 $t_p = 20$  ns

*Nutation*

$$\omega_{\text{eff}} = \sqrt{\Omega^2 + \omega_1^2}$$

$$\theta = \text{atan}(\omega_1/\Omega)$$

precession about the effective field



- the Fourier transform of a boxcar function approximates the excitation profile for a *small* flip angle:

$$P(\Omega) \approx \frac{\sin(\Omega t_p/2)}{\Omega t_p/2}$$



# Overcoming the microwave power limit to bandwidth

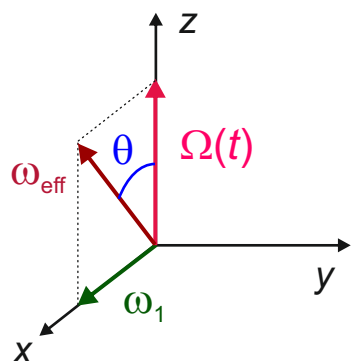
The power-bandwidth relation for rectangular monochromatic pulses

$$BW \approx 1/t_p \quad t_p = \beta/\omega_1 \quad \Rightarrow BW \approx \omega_1/\beta = g_x \mu_B B_1/(\hbar\beta) \quad \text{with } B_1 \propto \sqrt{P_{mw}}$$

! the time available for spin excitation is of the order of  $T_2$  (microseconds) or  $T_1$  (milliseconds), coherent excitation saturation  
 whereas spectral width is of the order of 100 MHz ... 10 GHz, corresponding to  $t_p = 10 \dots 0.1$  ns

## Frequency-swept pulses

Accelerated frame picture

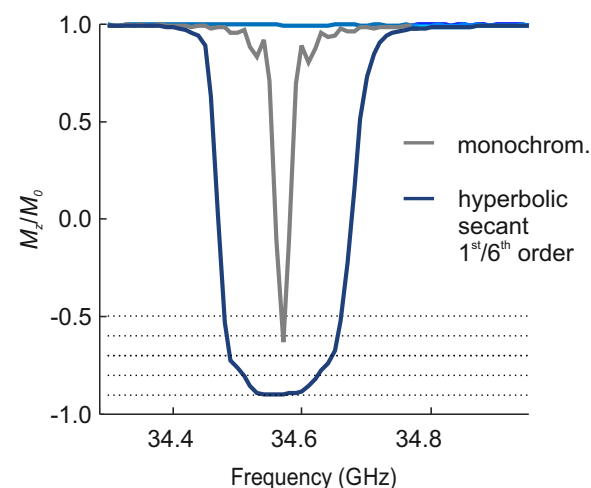


Adiabaticity

$$Q = \omega_{eff}/|d\theta/dt|$$

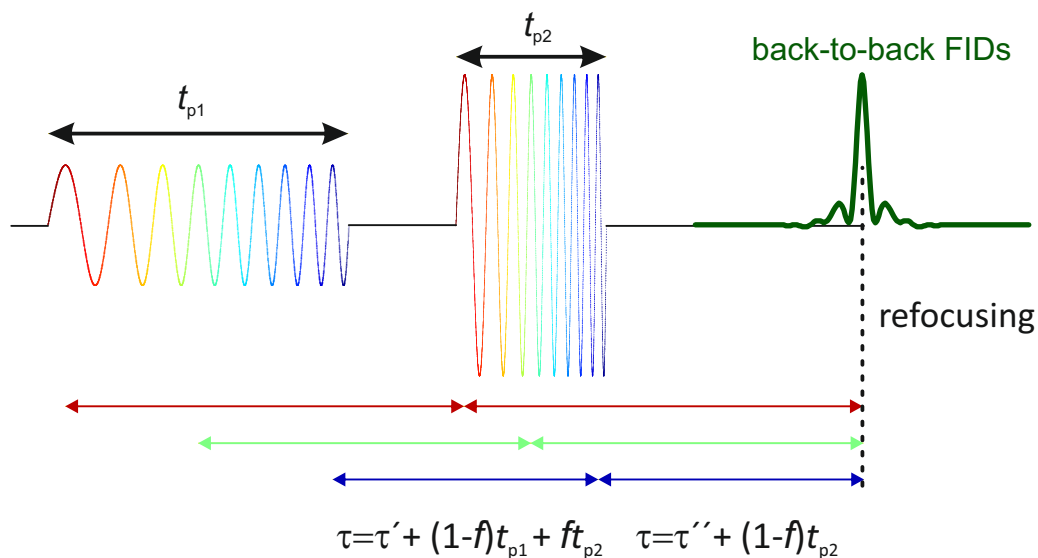
- for  $Q \gg 1$  the magnetization vector follows the effective field from +z to -z  
 $\Rightarrow$  **perfect** wideband inversion

Well, nobody is perfect



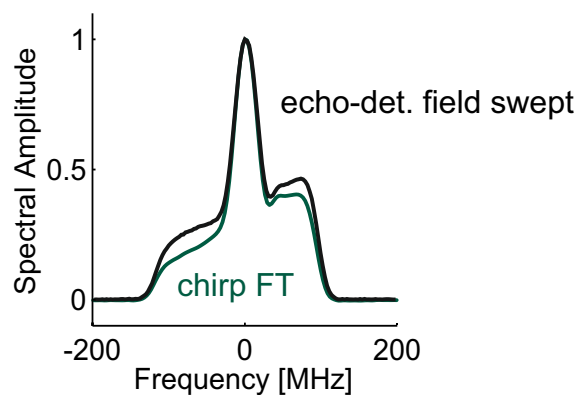
# Chirp echoes

## Primary chirp echo

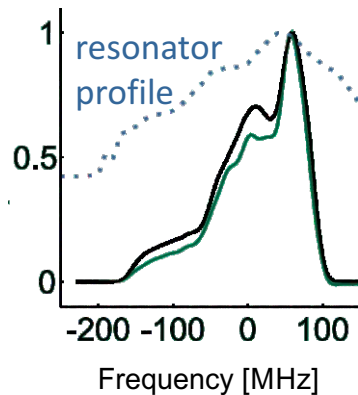


- a spin packet  $\Omega$  is excited after time fraction  $f = (\Omega - \omega_{\text{start}})/B$
- for all packets to refocus at the same time, we must have  $\tau' + (1-f)t_{p1} + ft_{p2} = \tau'' + (1-f)t_{p2}$  for all  $f$ , which gives  $t_{p1} = 2t_{p2}$
- Fourier transformation of the echo gives the frequency-domain EPR spectrum

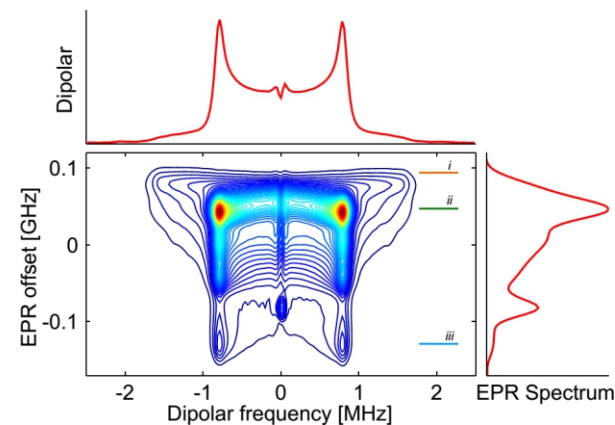
## X-band chirp FT EPR



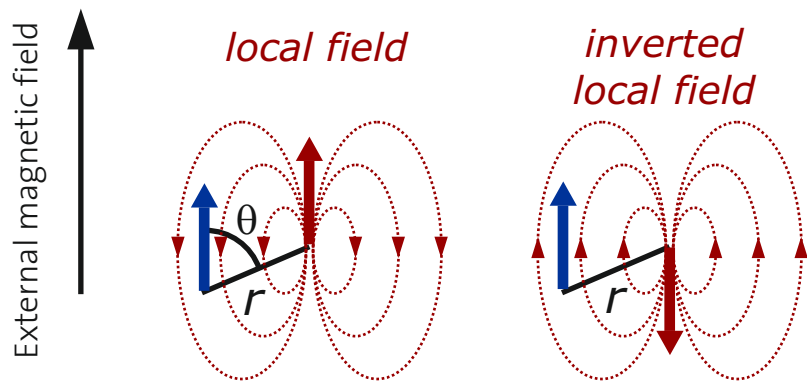
## Q-band chirp FT EPR



## Q-band SIFTER/FT EPR correlation



# Dipole-dipole coupling



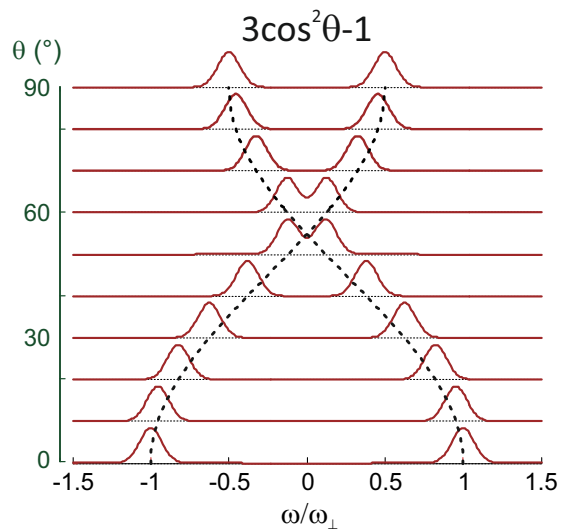
spin vector  
operators

$$\hat{H}_{dd} = \frac{1}{r_{SI}^3} \frac{\mu_0 \hbar}{4\pi} \gamma_S \gamma_I \left[ \hat{\mathbf{S}} \cdot \hat{\mathbf{I}} - 3 \frac{1}{r_{SI}^2} (\hat{\mathbf{S}} \cdot \mathbf{r}_{SI}) (\hat{\mathbf{I}} \cdot \mathbf{r}_{SI}) \right]$$

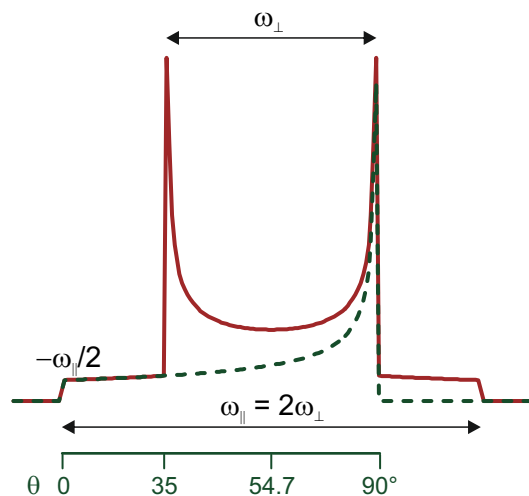
Relevant terms

$$\frac{1}{r_{SI}^3} \frac{\mu_0 \hbar}{4\pi} \gamma_S \gamma_I (3\cos^2\theta - 1) [\hat{S}_z \hat{I}_z - \frac{1}{2}(\hat{S}^+ \hat{I}^- + \hat{S}^- \hat{I}^+)]$$

Orientation dependence



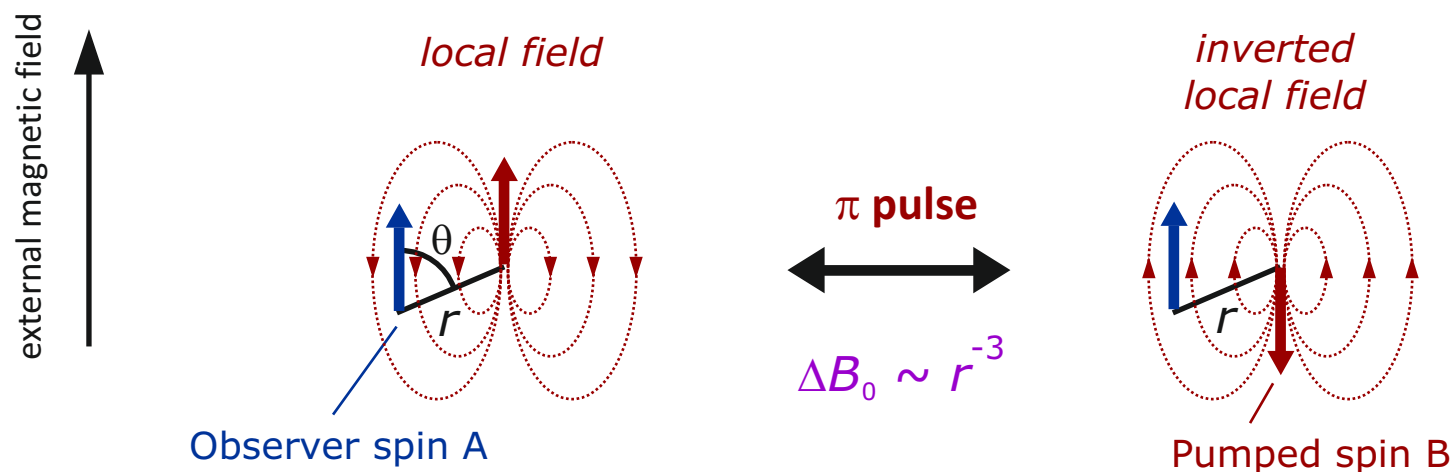
Pake pattern



To infer the distance  $r_{SI}$ , we must

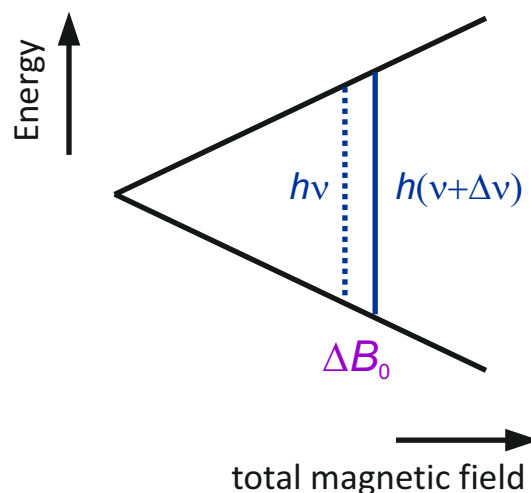
- separate dipole-dipole coupling from stronger interactions
- get rid of the orientation dependence

# Separation of the electron-electron coupling from other interactions



1. Signal is labeled with resonance frequency  $\nu = g\mu_B B_0 / h$

2. Pumped spin is inverted, field shifts



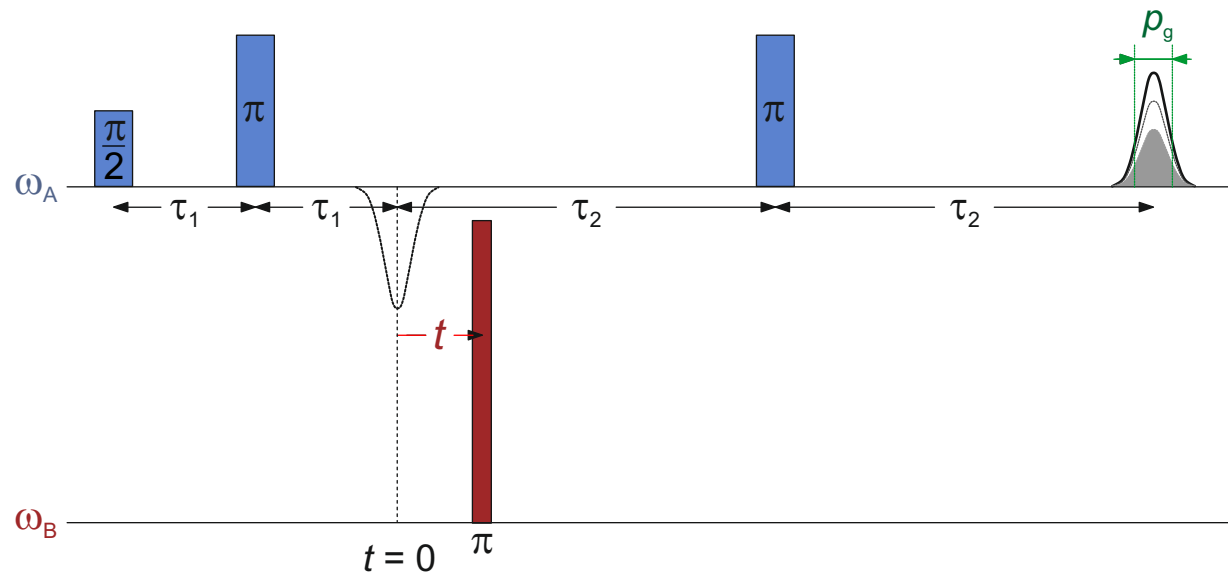
3. Field shift  $\Delta B_0$  is measured via change  $\Delta\nu$  of resonance frequency

4. Distance  $r$  is computed from  $\Delta\nu$

$$r^3 = (52.04 \text{ MHz}/\Delta\nu) \text{ nm}$$

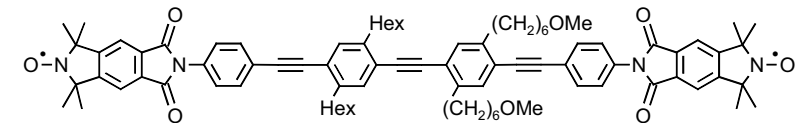
# What is DEER?

a proud animal that also accepts to be named PELDOR

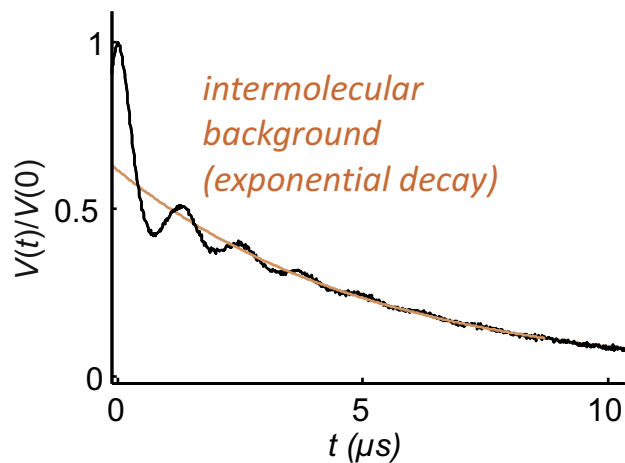


The **echo amplitude** is observed as a function of **time t**

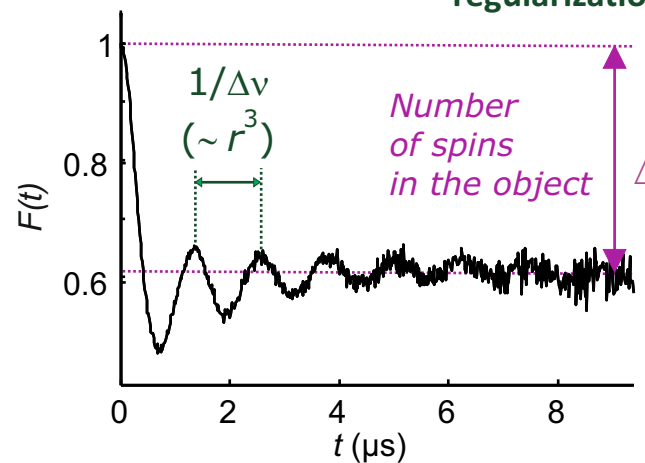
## Model compound



## Primary data

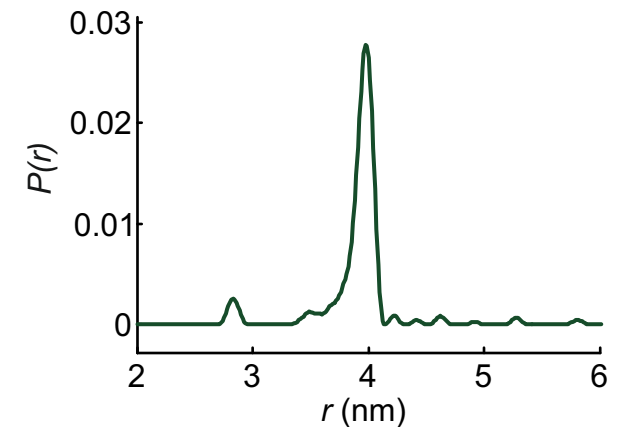


## Form factor



Tikhonov regularization

## Distance distribution



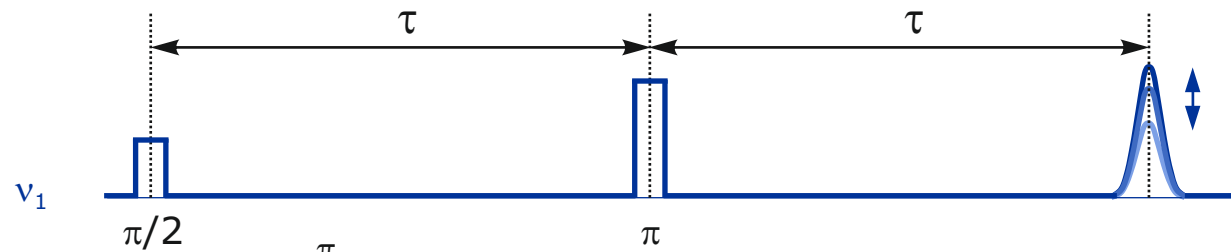
MARTIN RE et al., *Angew. Chem. Int. Ed.* **1998**, *37*, 2834

PANNIER M, VEIT S, GODT A, JESCHKE G, SPIESS HW, *J. Magn. Reson.* **2000**, *142*, 331

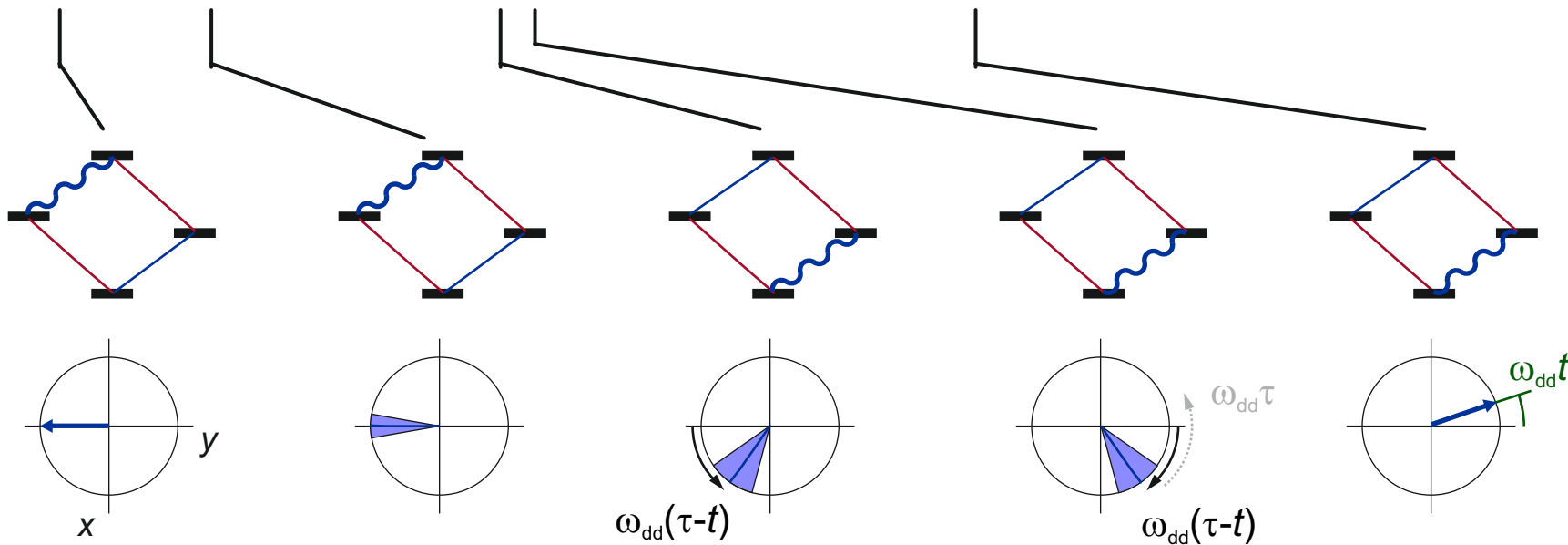
JESCHKE G et al. *J. Magn. Reson.* **2002**, *155*, 72

JESCHKE G et al. *Appl. Magn. Reson.* **2006**, *30*, 473

# Separating the dipole-dipole interaction by a spin flip



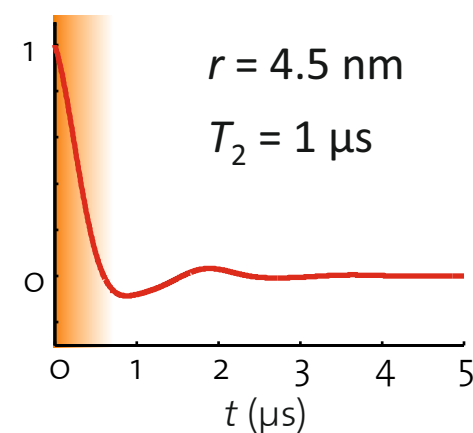
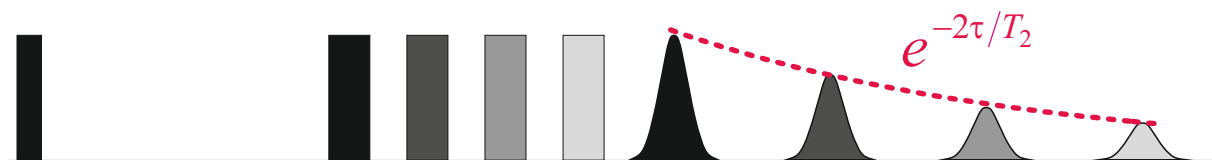
- constant delay  $\tau$ :  
no echo variation  
due to relaxation
- variable delay  $t$ :  
resonance frequency  
is changed at different times



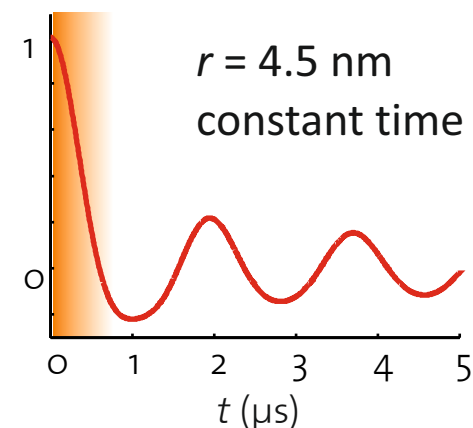
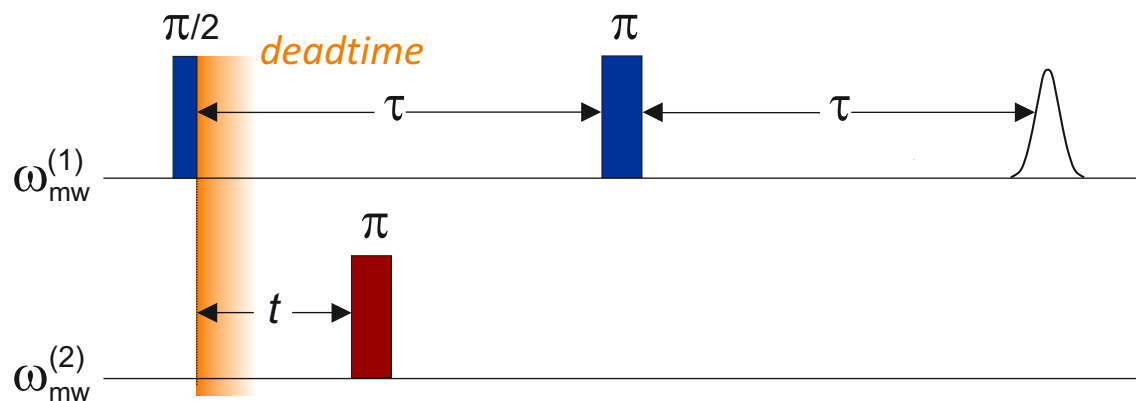
original experiment by Salikhov, Milov, Shirov, and Tsvetkov

## Ultimate resolution by a constant-time experiment

- if the total length of the pulse sequence changes, dipolar evolution is damped by relaxation
- for distances larger than ~3-4 nm, dipolar modulation is overdamped

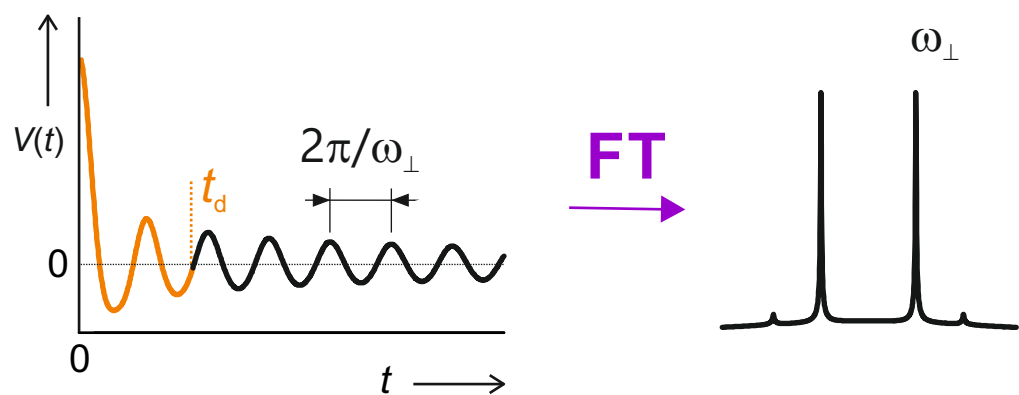


Keep  $\tau$  fixed, vary only  $t$ , pulse sequence has constant length



We trade sensitivity for resolution

# Eliminating dead-time by the four-pulse DEER experiment



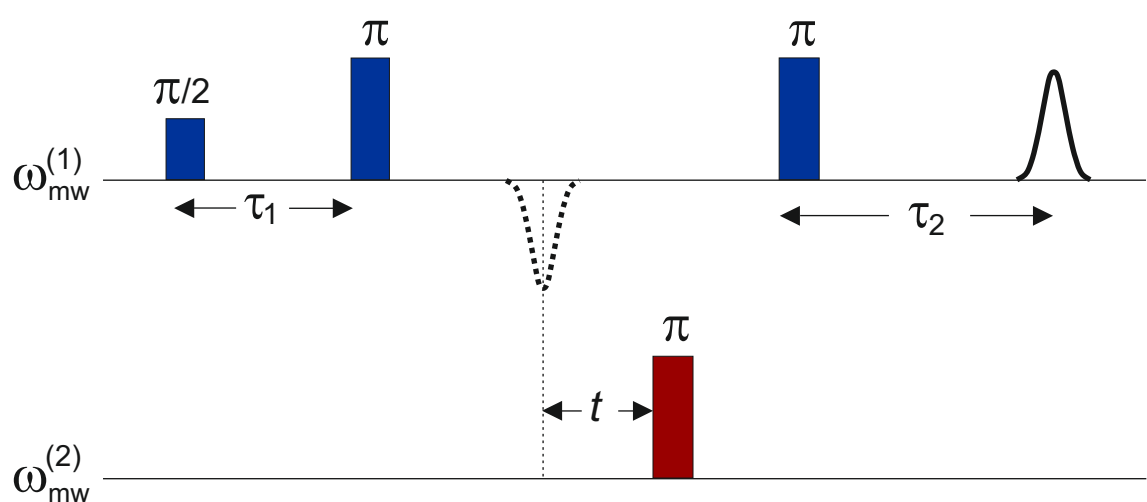
unreliable results due to deadtime for

- distributions with  $r < 2.2$  nm
- direct extraction of distance distributions

sensitivity loss for small  $r$

Remedy: Refocusing of the echo

$$\Rightarrow \text{accessible } t: -\tau_1 + t_d < t < \tau_2 - t_d$$



Typical values (nitroxides):

$\tau_1 \approx 120 \dots 200$  ns, 400 ns with deuterated matrix

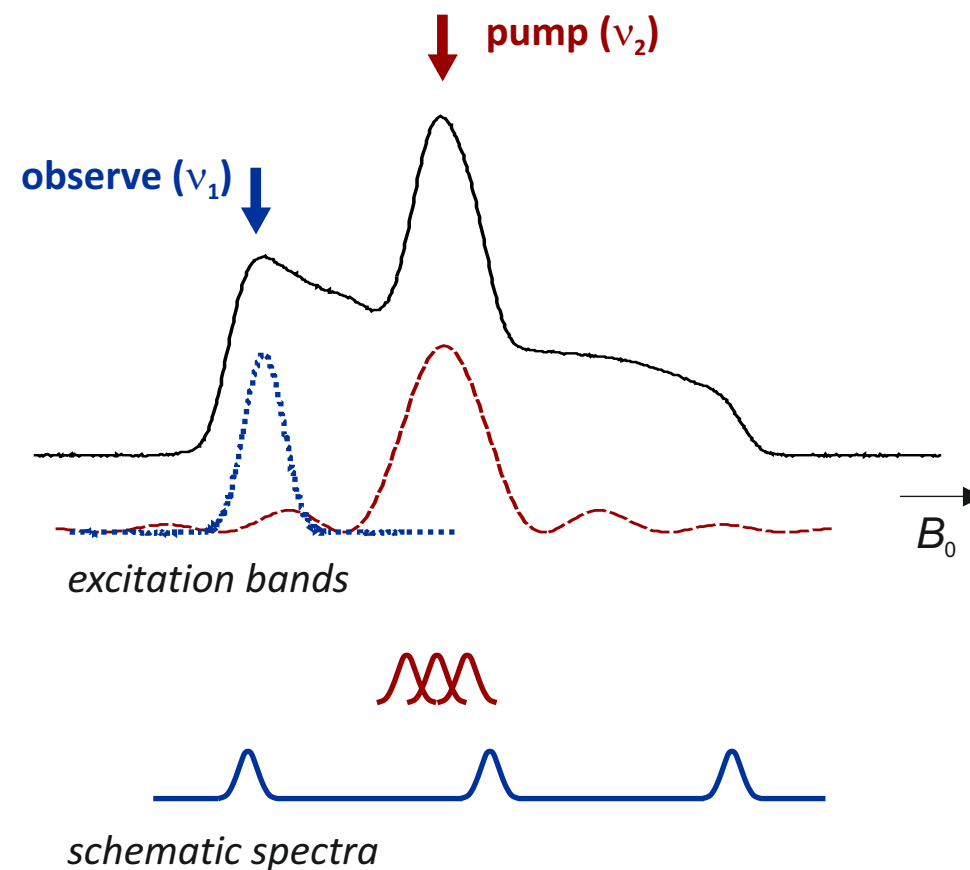
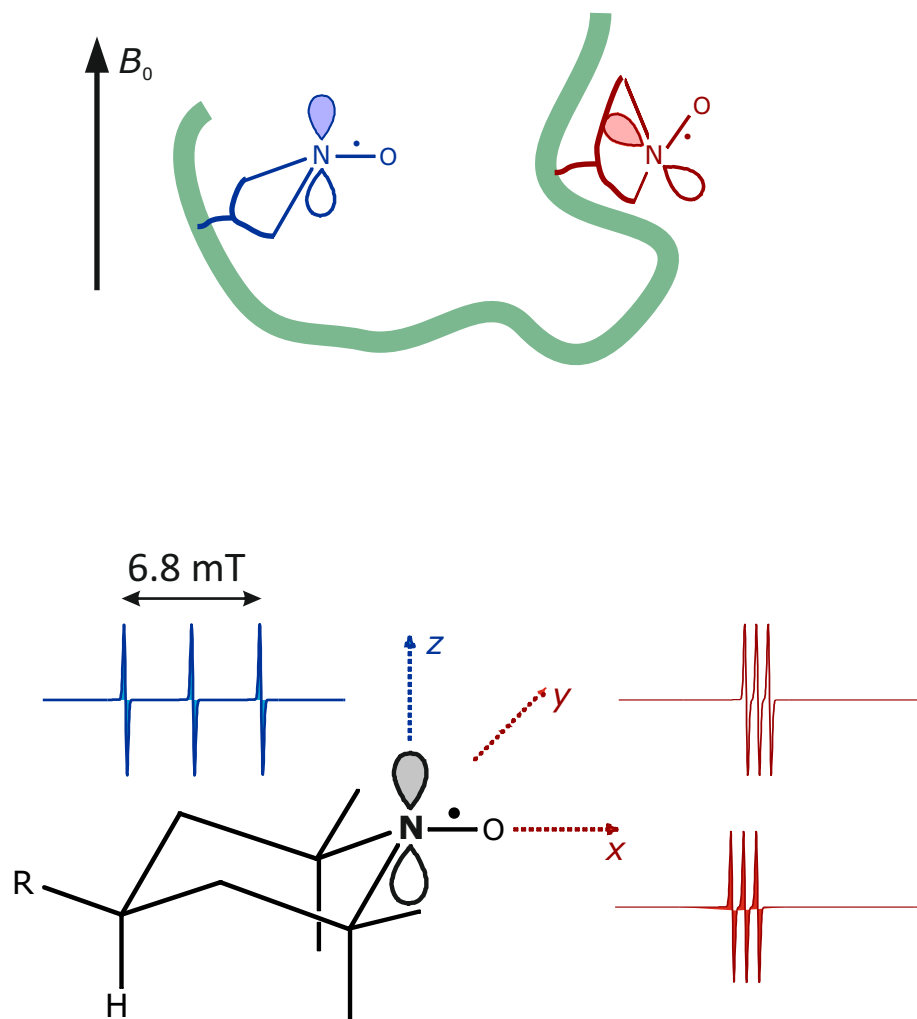
$\tau_2 \approx 1 \mu\text{s}$  (2 nm) ... 6  $\mu\text{s}$  (5-8 nm)

$t_p(1) = 32$  ns

$t_p(2) = 12$  ns (opt. sensitivity)



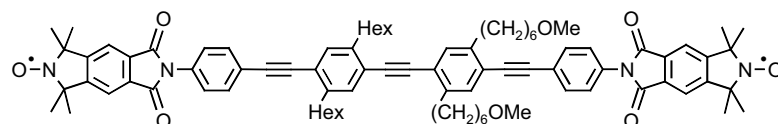
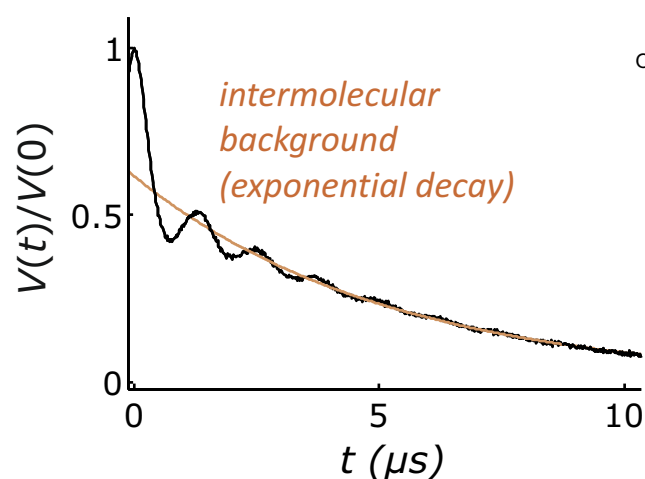
## Why can we excite two nitroxide labels with two frequencies?



- only for a fraction  $\lambda$  of all spin pairs of the **observed spin A** the **coupled spin B** is indeed excited (pumped)

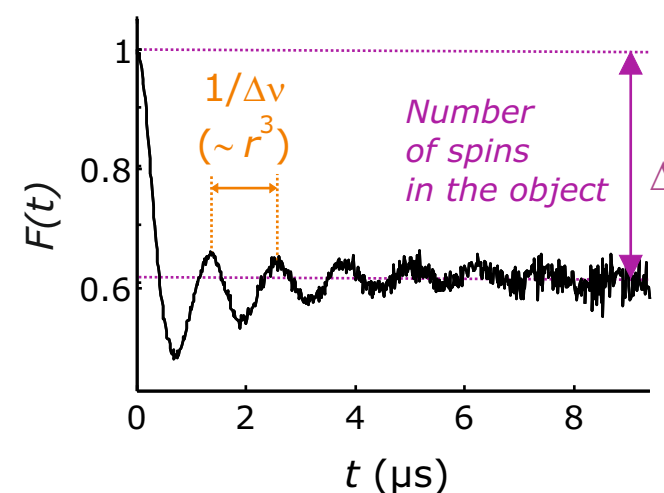
## Information content of the DEER signal

### Primary experimental data



background  
correction

### Form factor (intramolecular)



- decay envelope depends on distance distribution: fast decay of modulation  $\Rightarrow$  broad distribution

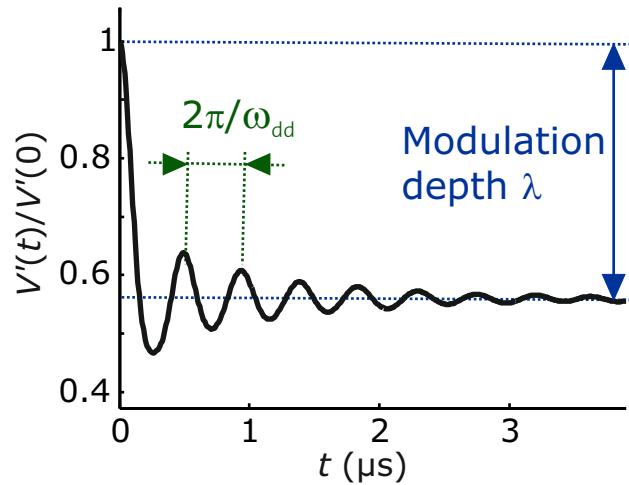
Dipole-dipole interaction 
$$\omega_{dd} = \frac{1}{r^3} \frac{\mu_0}{4\pi h} g_1 g_2 \mu_B^2$$

Dipolar frequency 
$$d(\theta) = (1 - 3 \cos^2 \theta) \omega_{dd}$$

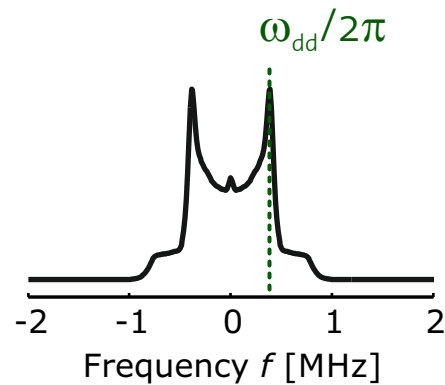
Form factor of an isolated spin pair 
$$F(r, t) = \frac{V(t)}{V(0)} = 1 - \int_0^{\pi/2} \lambda(\theta) \{1 - \cos[2\pi\nu_{dd}(\theta)]\} \sin\theta d\theta$$

# Data analysis by Fourier transformation

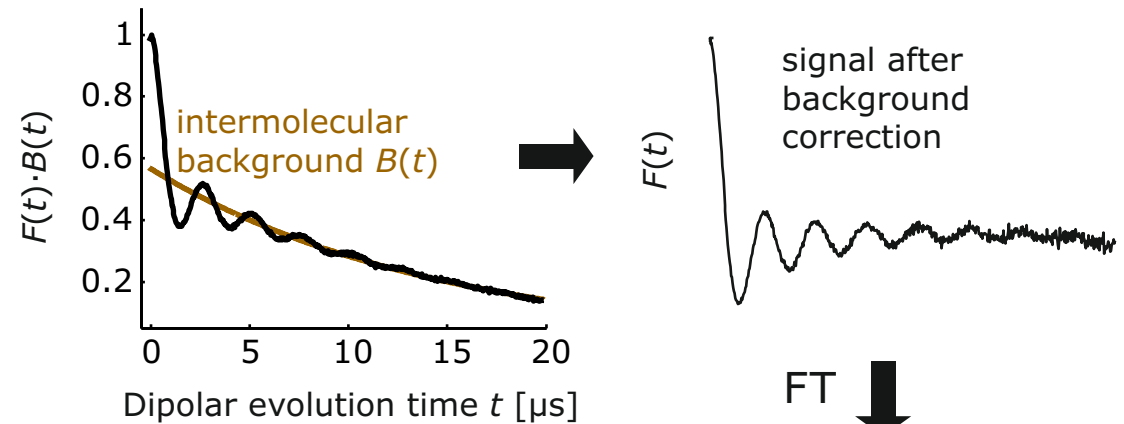
## Dipolar evolution function (form factor)



Fourier transformation



## Normalized echo modulation

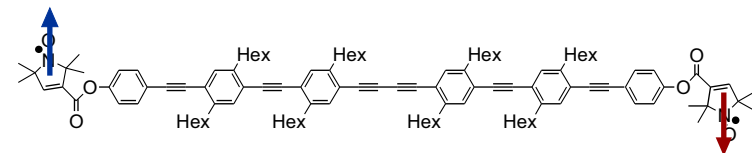
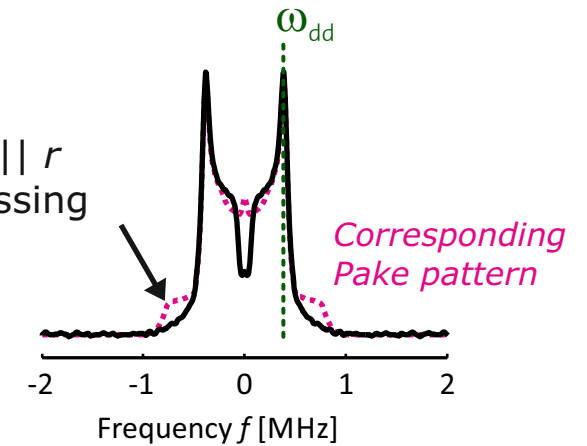


FT

$$\omega_{dd}/2\pi = 380 \text{ kHz}$$

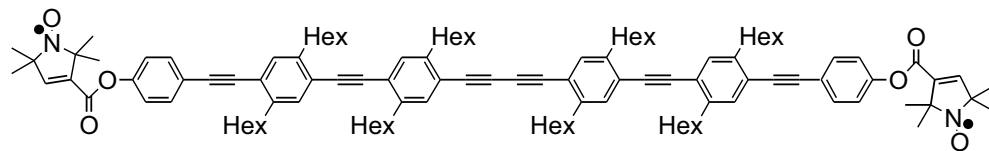
$$\Rightarrow r = 5.05 \text{ nm}$$

$B_0 \parallel r$   
missing

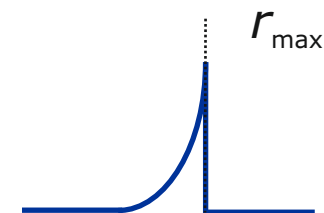


# Average distance is not the full information

Molecular dynamics causes distribution of  $r$

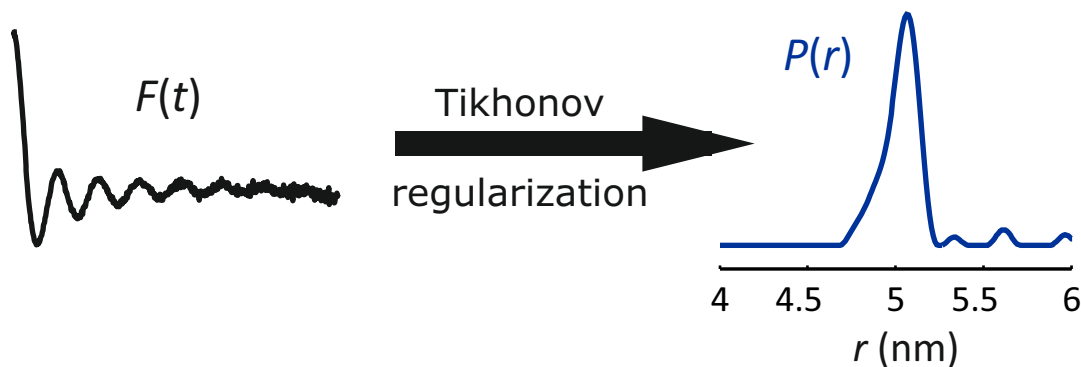


Expectation:



bent, higher energy, lower population, smaller  $r$   
 stretched, min. energy, max. population, max.  $r$

Experimental distance distribution



Basic mathematics

fit simulated dipolar evolution function  $S(t) = K P(r)$

Minimize

$$G_{\alpha}(P) = \| KP(r) - F(t) \|^2 + \alpha \left\| \frac{d^2}{dr^2} P(r) \right\|^2$$

mean square deviation

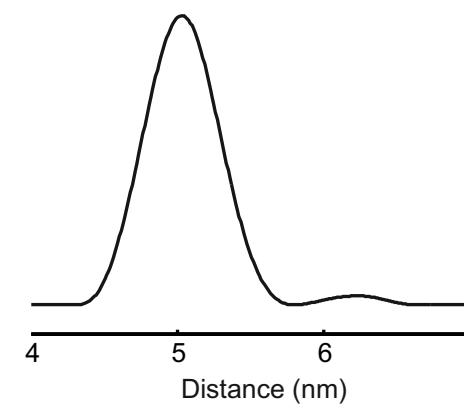
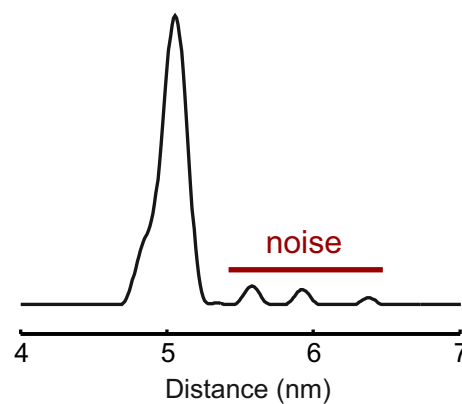
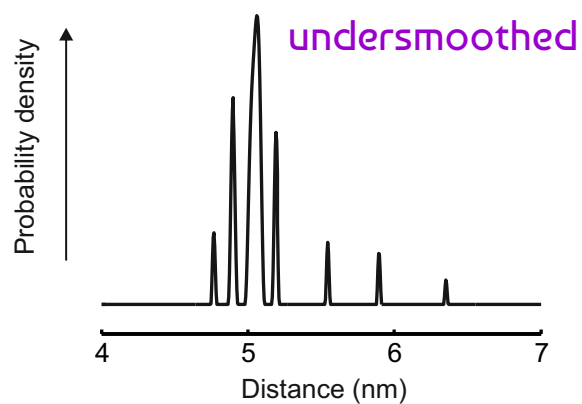
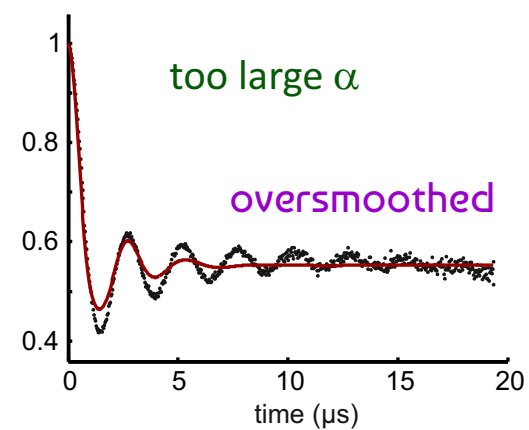
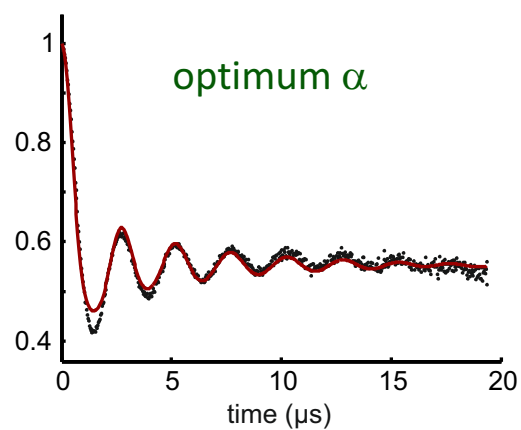
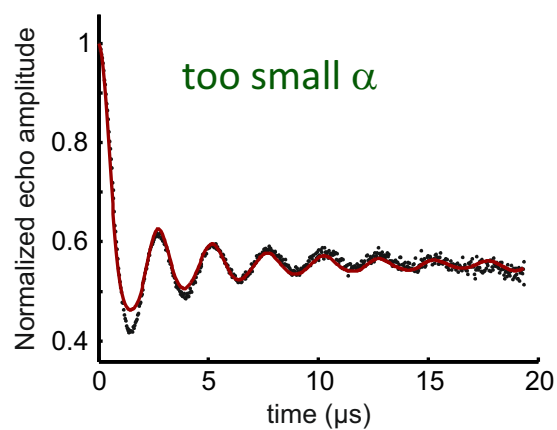
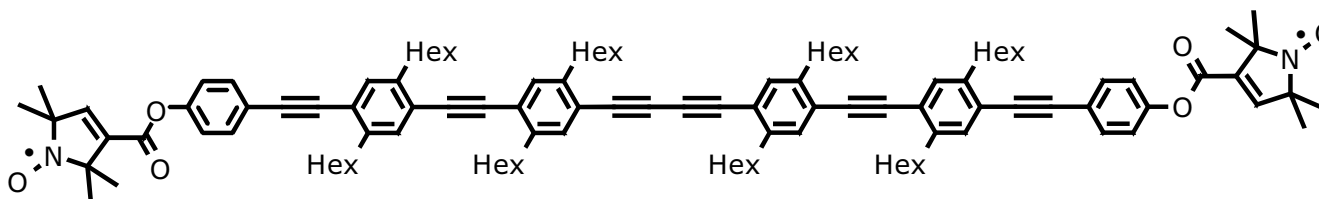
roughness of 2<sup>nd</sup> derivative

regularization parameter  $\alpha$

G. JESCHKE, A. KOCH, U. JONAS, A. GODT A, *J. Magn. Reson.* 155, 72-82 (2002)

G. JESCHKE et al. *Appl. Magn. Reson.* 30, 473-498 (2006)

## Why regularization is required and how to do it

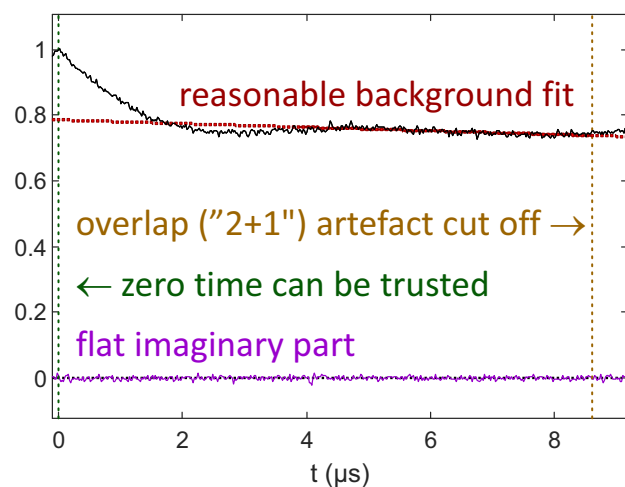


# Don't forget to validate your data

## Don't forget to validate your data

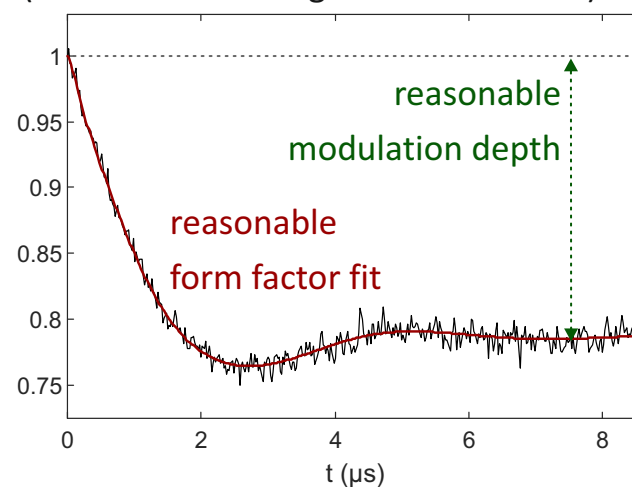
PTBP1/EMCV-IRES(D-F) 85 kDa protein-RNA complex, protein double mutant 235R1/388R1

### Primary data

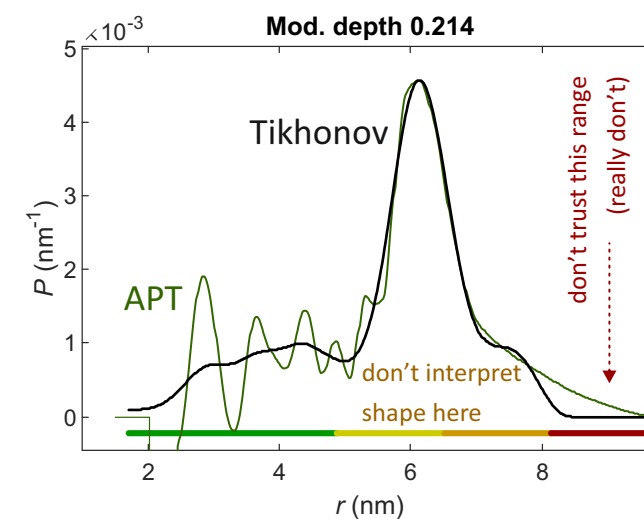


### Form factor

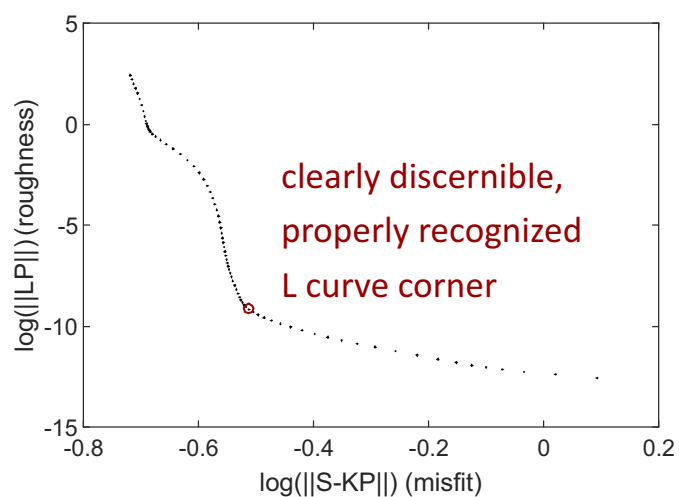
(with reasonable signal-to-noise ratio)



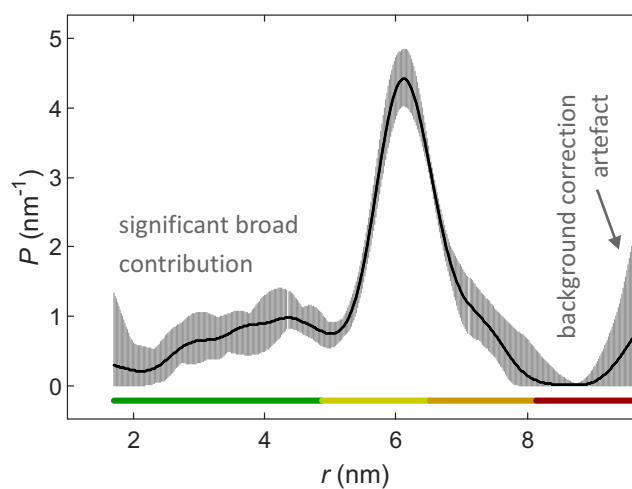
### Distance distribution



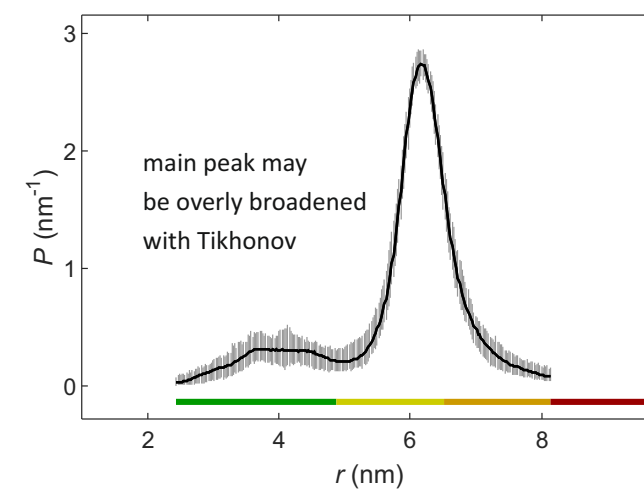
### L curve



### Validation



### Second opinion (DEERNet)



## Further reading & useful stuff on pulsed EPR

### Product operator formalism and phase cycling

R.R. ERNST, G. BODENHAUSEN, A. WOKAUN, *Principles of Nuclear Magnetic Resonance in One and Two Dimensions*, Oxford University Press, Oxford, 1987.

O. W. SØRENSEN, *Progr. Nucl. Magn. Reson. Spectrosc.* 21, 503-570 (1989).

### Topics specific to pulse EPR

A. SCHWEIGER, G. JESCHKE, *Principles of Pulse Electron Paramagnetic Resonance*, Oxford University Press, Oxford, 2001.

### Adiabatic & fast passage

J. BAUM, R. TYCKO, A. PINES, *Phys.Rev. A* 32, 3435–3447 (1985).

A. DOLL, G. JESCHKE, *J. Magn. Reson.* 280, 46-62 (2017)

### Numerical computations

EasySpin, Spinach

### Product operator computations & beyond

SpinDynamica

### Signal processing

G. JESCHKE, *Lecture notes, ETH Zürich*

[epr.ethz.ch/education/messtechnik.html](http://epr.ethz.ch/education/messtechnik.html)

### (Pulsed) EPR lecture notes

G. JESCHKE, *Lecture notes, ETH Zürich*

[epr.ethz.ch/education.html](http://epr.ethz.ch/education.html)

## Further reading & useful stuff on DEER & related

### General overview on techniques for distance distribution measurements

P. P. BORBAT, J. H. FREED, eMagRes, Wiley, Vol. 5, 465–494 (2017). *Dipolar Spectroscopy – Single-Resonance Methods*

G. JESCHKE, eMagRes, Wiley, Vol. 5, 1459-1475 (2016). *Dipolar Spectroscopy – Double-Resonance Methods*

### Biological applications

G. JESCHKE, Annu. Rev. Phys. Chem. 63, 419-446 (2012), *DEER distance measurements on proteins*

B. ENDEWARD, A. MARKO, V. P. DENYSENKOV, S. T. SIGURDSSON, T. F. PRISNER, Methods in Enzymology 564, 403-425 (2015) *Advanced EPR Methods for Studying Conformational Dynamics of Nucleic Acids*

G. JESCHKE, Emerging Topics in Life Science, 2, 9-18, (2018) *The contribution of modern EPR to structural biology*

### Structure modelling based on distance distributions

G. JESCHKE, Proteins 84, 544-560 (2016)

G. HAGELUEKEN, D. ABDULLIN, O. SCHIEMANN, Methods in Enzymology, 563, 595-622 (2015) *mtsslSuite: Probing Biomolecular Conformation by Spin-Labeling Studies*

G. JESCHKE, Protein. Sci. 27, 76-85 (2018)

*MMM - A toolbox for integrative structure modelling*

### DeerAnalysis & MMM

[epr.ethz.ch/software.html](http://epr.ethz.ch/software.html)

[deeranalysis.org](http://deeranalysis.org)

